

Thursday, August 31, 2017

Basic Review of time-independent E & M
(in vacuum)

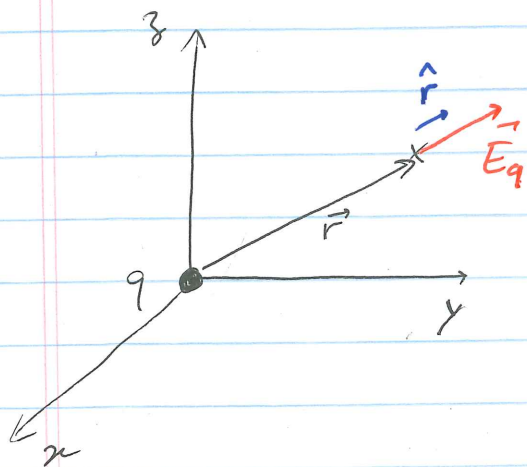
E-field satisfies: - $\nabla \cdot \vec{E} = \frac{1}{\epsilon_0} \rho(\vec{r})$ Gauss's Law

where $\epsilon_0 = 8.85 \times 10^{-12} \frac{C^2}{N \cdot m^2}$
= permittivity of free space

ρ = charge density in $\frac{\text{Coulombs (C)}}{m^3}$

- $\nabla \times \vec{E} = 0$

E-field of a point charge: $\vec{E}_q = \frac{1}{4\pi\epsilon_0} \frac{q}{|\vec{r}|^2} \hat{r}$



charge density of a point charge q :

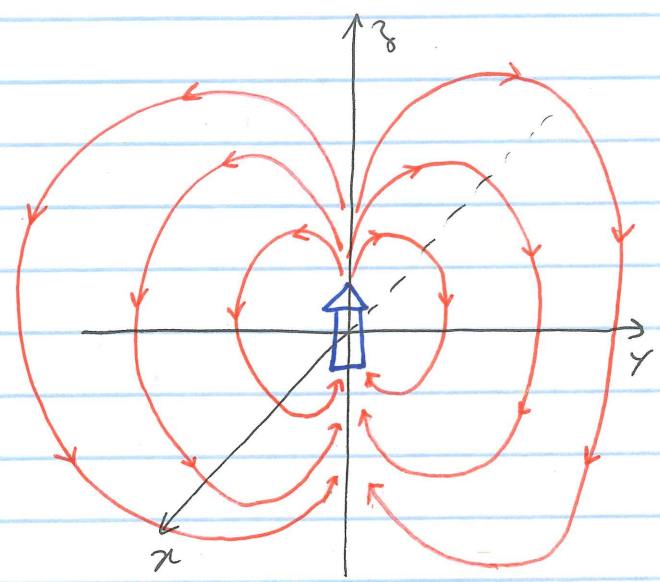
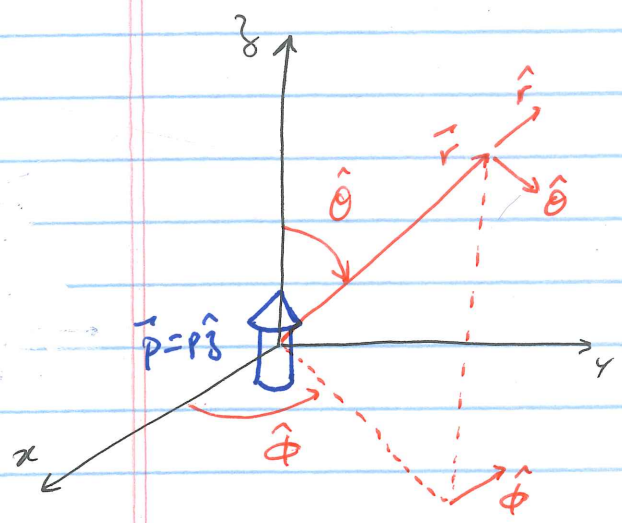
$$\rho(\vec{r}) = q \delta^3(\vec{r} - \vec{r}_0)$$

$$= q \delta(x - x_0) \delta(y - y_0) \delta(z - z_0)$$

Electric field of an ideal electric dipole \vec{P} :

$$\vec{E}_{\text{dipole}} = \frac{|\vec{P}|}{4\pi\epsilon_0} \frac{1}{r^3} \left(2 \cos\theta \hat{r} + \sin\theta \hat{\theta} \right)$$

$3(\vec{P} \cdot \hat{r})\hat{r} - \vec{P}$
 "Coordinate free" version



Cylindrically symmetric
(also spherical symmetry)

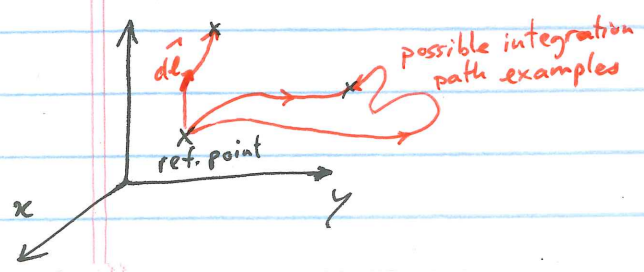
$$\vec{P} = \int \vec{r} \rho(\vec{r}) d^3r$$

$$= q \vec{d} \quad \text{for}$$



Potential formulation of Electrostatics (derived from $\nabla \times \vec{E} = 0$)

Electric potential: $V(\vec{r}) = - \int_{\text{ref. point}}^{\vec{r}} \vec{E}(\vec{r}') \cdot d\vec{l}'$



reference point frequently chosen at $\vec{r} = \text{infinity}$

does not depend on integration path

⚠ but does depend on the reference point

Potential difference does not depend on the ~~reference~~ reference point.

$$\begin{aligned}\Delta V &= V(\vec{r}_b) - V(\vec{r}_a) \\ &= - \int_{\text{ref. point}}^{\vec{r}_b} \vec{E} \cdot d\vec{l} - \left(- \int_{\text{ref. point}}^{\vec{r}_a} \vec{E} \cdot d\vec{l} \right) \\ &= - \int_{\text{ref. point}}^{\vec{r}_b} \vec{E} \cdot d\vec{l} - \int_{\vec{r}_a}^{\text{ref. point}} \vec{E} \cdot d\vec{l} = - \int_{\vec{r}_a}^{\vec{r}_b} \vec{E} \cdot d\vec{l}\end{aligned}$$

note: $\vec{\nabla} V = \left(\lim_{\Delta x \rightarrow 0} \frac{V(\vec{r} + \Delta x \hat{x}) - V(\vec{r})}{\Delta x}, \lim_{\Delta y \rightarrow 0} \frac{V(\vec{r} + \Delta y \hat{y}) - V(\vec{r})}{\Delta y}, \right.$

$\left. \lim_{\Delta z \rightarrow 0} \frac{V(\vec{r} + \Delta z \hat{z}) - V(\vec{r})}{\Delta z} \right)$

$$= \lim_{\Delta x, \Delta y, \Delta z \rightarrow 0} \left(\underbrace{-\frac{1}{\Delta x} \int_{\vec{r}}^{\vec{r} + \Delta x \hat{x}} \vec{E} \cdot d\vec{l}}_{-E_x}, \underbrace{-\frac{1}{\Delta y} \int_{\vec{r}}^{\vec{r} + \Delta y \hat{y}} \vec{E} \cdot d\vec{l}}_{-E_y}, \underbrace{-\frac{1}{\Delta z} \int_{\vec{r}}^{\vec{r} + \Delta z \hat{z}} \vec{E} \cdot d\vec{l}}_{-E_z} \right)$$

$$= -\vec{E}$$

$$\Rightarrow \boxed{\vec{E} = -\vec{\nabla} V} \quad \text{differential formulation}$$

Potential of a point charge: $V_q(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{q}{|\vec{r}|}$
(ref. point at infinity)

Potential of an ideal dipole: $V_{\text{dipole}}(\vec{r}) = \frac{\vec{p} \cdot \vec{r}}{4\pi\epsilon_0 r^2} = \frac{|\vec{p}| \cos\theta}{4\pi\epsilon_0 r^2}$

$$\left. \begin{aligned} \vec{\nabla} \cdot \vec{E} &= \frac{\rho(\vec{r})}{\epsilon_0} \\ \vec{\nabla} \times \vec{E} &= 0 \\ (\text{i.e. } \vec{E} &= -\vec{\nabla} V) \end{aligned} \right\} \Rightarrow \boxed{\vec{\nabla}^2 V = -\frac{\rho}{\epsilon_0}} \quad \begin{array}{l} \text{Poisson's} \\ \text{equation} \end{array}$$

$$\uparrow$$

$$\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Energy: Electrostatic potential energy ($\Delta U = q \Delta V$)

$$U = \frac{\epsilon_0}{2} \int_{\text{all space}} |\vec{E}|^2 d^3r$$

The magnetic field satisfies:

$$\boxed{\vec{\nabla} \cdot \vec{B} = 0}$$

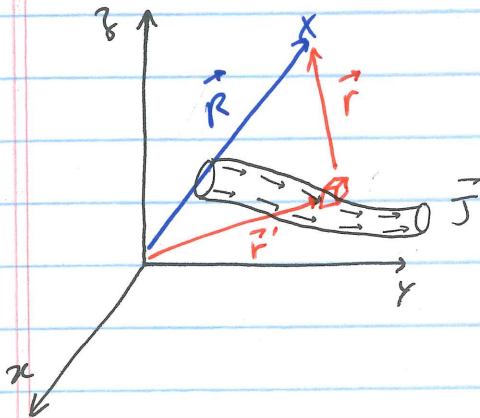
"no magnetic monopoles" law

$$\boxed{\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}}$$

Ampère's Law

Biot-Savart Law:
$$\vec{B}(\vec{R}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}') \times \hat{r}}{r^2} d^3r'$$

$\hat{r} = \frac{\vec{R} - \vec{r}'}{|\vec{R} - \vec{r}'|}$



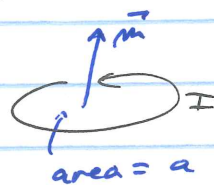
where $\mu_0 = 4\pi \times 10^{-7} \text{ N/A}^2$
= permeability of free space

$\vec{J}(\vec{r})$ = current density
(A/m²)

Magnetic field of a magnetic dipole: $\vec{m} = I \vec{a}$

$$\vec{B}_{\text{dipole}} = \frac{\mu_0 |\vec{m}|}{4\pi r^3} \left(2 \cos\theta \hat{r} + \sin\theta \hat{\theta} \right)$$

$$3(\vec{m} \cdot \hat{r})\hat{r} - \vec{m}$$



↳ same form as \vec{E} -field of an electric dipole!

Magnetic Vector Potential \vec{A} :

$$\vec{B} = \nabla \times \vec{A}$$

(definition)

does not fully specify \vec{A}

$$\nabla^2 \vec{A} = -\mu_0 \vec{J}$$

potential formulation of Ampère's law

if we choose \vec{A} such that $\nabla \cdot \vec{A} = 0$

$$\vec{A}(\vec{R}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{|\vec{r}|} d^3r'$$

for a magnetic dipole: $\vec{A}_{\text{dipole}} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \hat{r}}{r^2}$

Magnetostatic Energy: $U = \frac{1}{2} \frac{1}{\mu_0} \int_{\text{all space}} |\vec{B}|^2 d^3r$

(we will see this next week)

Lorentz Force Law:

$$\vec{F} = \text{force on a charge } q$$

$$= q\vec{E} + q\vec{v} \times \vec{B}$$

(does not change when \vec{E} & \vec{B} are time-dependent)

Fundamental Vector Integral Theorems

$$\int_{\vec{r}_a}^{\vec{r}_b} (\vec{\nabla} f) \cdot d\vec{\ell} = f(\vec{r}_b) - f(\vec{r}_a) \quad \text{"gradient theorem"}$$

$$\int_{\text{Volume}} (\underbrace{\vec{\nabla} \cdot \vec{F}}_{\text{divergence}}) d^3r = \oint_{\text{bounding surface}} \vec{F} \cdot d\vec{s} \quad \begin{array}{l} \text{Gauss's / Green's} \\ \text{Theorem} \\ \text{or divergence theorem} \end{array}$$

$$\int_{\text{Surface}} (\underbrace{\vec{\nabla} \times \vec{F}}_{\text{curl}}) \cdot d\vec{s} = \oint_{\text{bounding line}} \vec{F} \cdot d\vec{\ell} \quad \text{Stoke's theorem}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} \Leftrightarrow \oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enclosed}} \quad \text{Ampère's law}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \Leftrightarrow \oint_{\text{bounding surface}} \vec{E} \cdot d\vec{s} = \frac{q_{\text{enclosed}}}{\epsilon_0} \quad \text{Gauss's law}$$



Tuesday quiz (5-10 minutes)