

Thursday, October 26, 2017

#1

Reflection & Transmission at a conducting surface

Recall boundary conditions:

1) $\epsilon_1 E_{1,\perp} - \epsilon_2 E_{2,\perp} = \sigma_f = 0$

2) $\vec{E}_{1,\parallel} = \vec{E}_{2,\parallel}$

3) $B_{1,\perp} = B_{2,\perp}$

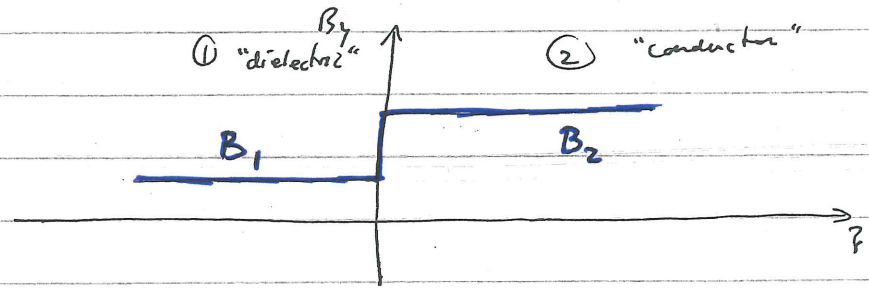
4) $\frac{1}{\mu_1} \vec{B}_{1,\parallel} - \frac{1}{\mu_2} \vec{B}_{2,\parallel} = \vec{K}_f \times \hat{n}$

$\vec{K}_f = 0$

(i.e. no free surface current ... however $\vec{J}_f \neq 0$)

↳ why? ... suppose $\mu_1 \approx \mu_2 \approx \mu_0$

if $\vec{K}_f \neq 0$, then \vec{B} is discontinuous across the boundary



$\Rightarrow \frac{\partial B_y}{\partial z} = \underbrace{\beta \delta(z)}_{\text{surface current}} \Rightarrow \vec{E}_{\text{surface}} = +\infty$ (if you believe Ohm's law)

Ampere's law: $\nabla \times \vec{B} = \mu_0 \sigma \vec{E} + \mu_0 \epsilon \frac{\partial \vec{E}}{\partial t}$ focus on this term

$(\cancel{\partial_y B_z} - \cancel{\partial_z B_y}, \cancel{\partial_z B_x} - \cancel{\partial_x B_z}, \cancel{\partial_x B_y} - \cancel{\partial_y B_x}) = \left(\underbrace{\mu_0 \sigma + \mu_0 \epsilon \frac{\partial}{\partial t}}_{\text{leave out}} \right) (E_x, E_y, E_z)$
 $\beta \delta(z)$

$\Rightarrow E_x = +\infty$ (assuming σ is finite)

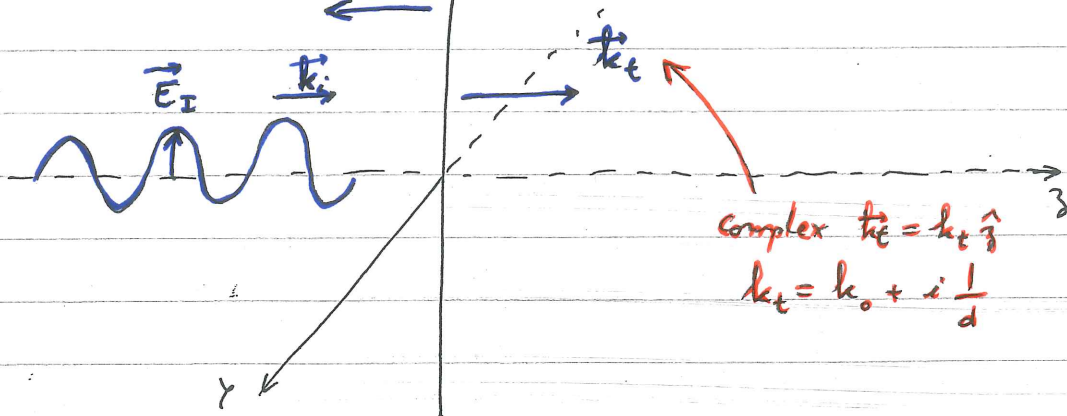
\Rightarrow also $\vec{K} \neq 0$ implies a surface $E_{\parallel} \rightarrow$ potential problem with b.c.

Normal incidence

dielectric n
 $\sigma = 0$

$k_r = -k_i$

good conductor
 $\sigma \gg \epsilon \omega$



complex $k_t = k_t \hat{z}$
 $k_t = k_0 + i \frac{1}{d}$

$$\vec{E}_I = E_0 e^{i(k_z z - \omega t)} \hat{x} \quad \text{and} \quad \vec{B}_I = \frac{n E_0}{c} e^{i(k_z z - \omega t)} \hat{y}$$
$$\vec{E}_R = E_{0,R} e^{i(-k_z z - \omega t)} \hat{x} \quad \text{and} \quad \vec{B}_R = -\frac{n E_{0,R}}{c} e^{i(-k_z z - \omega t)} \hat{y}$$
$$\vec{E}_T = E_{0,t} e^{i(k_t z - \omega t)} \hat{x} \quad \text{and} \quad \vec{B}_T = \frac{n(\omega)}{c} E_{0,t} e^{i(k_t z - \omega t)} \hat{y}$$

Complex $k_t = k_0 + i \frac{1}{d}$

$\frac{k_t E_{0,t}}{\omega}$

Boundary conditions

- 1) $E_{1,\perp} = 0 = E_{2,\perp}$
- 2) $\vec{E}_{1,\parallel} = \vec{E}_{2,\parallel} \Rightarrow E_0 + E_{0,R} = E_{0,t}$
- 3) $B_{1,\perp} = 0 = B_{2,\perp}$
- 4) $\vec{B}_{1,\parallel} = \vec{B}_{2,\parallel} \Rightarrow \frac{n E_0}{c} - \frac{n E_{0,R}}{c} = \frac{n(\omega)}{c} E_{0,t}$

$$\Leftrightarrow \frac{\hbar}{\omega} (E_{0,i} - E_{0,r}) = \frac{\hbar_t}{\omega} E_{0,t} \quad \leftarrow \text{Complex}$$

$$\Leftrightarrow E_0 - E_{0,r} = \frac{\hbar_t}{\hbar} E_{0,t} \quad (4)$$

$$(2) + (4) \Rightarrow 2E_0 = \left(1 + \frac{\hbar_t}{\hbar}\right) E_{0,t}$$

$$\Leftrightarrow t = \frac{E_{0,t}}{E_0} = \frac{2}{1 + \frac{\hbar_t}{\hbar}}$$

Complex

plug back into (2) $\Rightarrow E_{0,r} = E_{0,t} - E_0 = \left(\frac{2}{1 + \frac{\hbar_t}{\hbar}} - 1\right) E_0$

$$\Leftrightarrow r = \frac{E_{0,r}}{E_0} = \frac{1 - \hbar_t/\hbar}{1 + \hbar_t/\hbar}$$

Complex

$t \neq r \in \mathbb{C} \Rightarrow$ there are non-trivial phases between \vec{E}_I , \vec{E}_R , and \vec{E}_t

Perfect Conductor Limit: $\sigma \rightarrow +\infty$

In this limit $\hbar_t = \hbar_0 + i \frac{1}{d} \rightarrow +\infty$
($\sigma \rightarrow +\infty$)

$$= \frac{n\omega}{c} \sqrt{\frac{1}{2} \left(1 + \sqrt{1 + \left(\frac{\sigma}{\epsilon\omega}\right)^2}\right)} + i\omega \sqrt{\frac{\mu\epsilon}{2}} \sqrt{\sqrt{1 + \left(\frac{\sigma}{\epsilon\omega}\right)^2} - 1}$$

$$\rightarrow +\infty + i\infty$$

thus $\lim_{\sigma \rightarrow +\infty} t = 0$ & $\lim_{\sigma \rightarrow +\infty} r = -1$

$\Rightarrow E_{o,t} = 0 \Rightarrow$ b.c. 2 becomes $\vec{E}_{r,\parallel} = 0$

$$\Leftrightarrow \vec{E}_o + \vec{E}_{o,r} = 0$$

$$\Leftrightarrow \vec{E}_o = -\vec{E}_{o,r}$$

Note: In the $\sigma \rightarrow +\infty$ limit, you can get a surface current since $\vec{B}_{t,\parallel} = 0 \rightarrow$ b.c. 4 $\Rightarrow \frac{B_o \hat{y}}{\mu} - \left(-\frac{B_o \hat{y}}{\mu}\right) = \vec{K}_f \times \hat{z}$

$$\Rightarrow \left(\frac{\mu E_o}{c\mu} + \frac{\mu E_o}{c\mu} \right) \hat{z} = \vec{K}_f$$

$$\Leftrightarrow \vec{K}_f = \frac{2\mu \vec{E}_o}{c\mu}$$

\vec{E}_o does not diverge, because $\sigma \rightarrow +\infty$

Boundary conditions for a perfect conductor:

$$\vec{E}_{\parallel} = 0, \vec{B}_{\perp} = 0$$

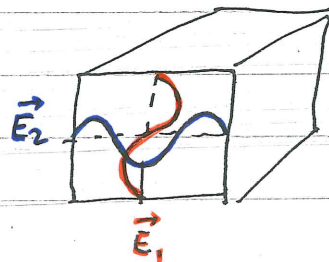
same as DC case.

Cavities (or Resonators):

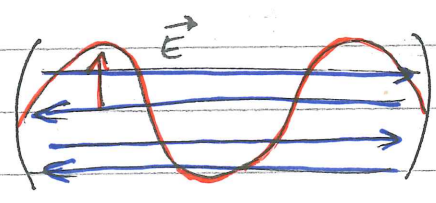
Definition: A container for EM waves.

ex1: A metal box

\hookrightarrow similar to a particle in a box (QM)



Ex 2: 2 mirrors, i.e. a Fabry-Perot cavity.



Most cavities have a part that is not fully reflective, but that is slightly transmissive.

example: a laser cavity

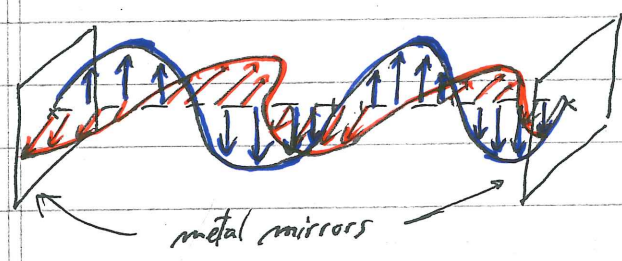
↳ light can get out on one side.

counter-example: Superconducting microwave cavity used by the S. Haroche group (ENS Paris) for 2012 Nobel prize work.

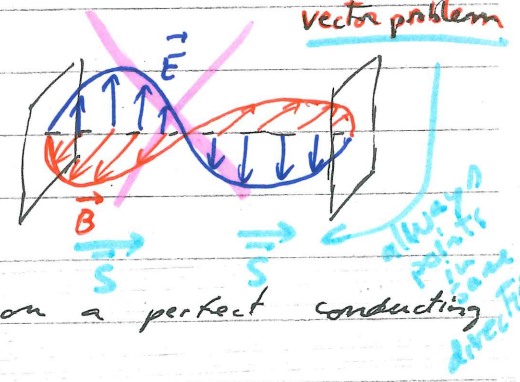
↳ How do they get the microwave photons in?

↳ black body radiation! (or atom beam)

boundary conditions (perfect conductor)



$$\begin{aligned} \vec{E}_{\parallel, \text{mirror}} &= 0 \\ \vec{B}_{\perp, \text{mirror}} &= 0 \end{aligned} \quad \left| \begin{array}{l} ? \\ \vec{B}_{\parallel, \text{mirror}} = 0 \end{array} \right. \begin{array}{l} \text{pointing} \\ \text{vector problem} \end{array}$$

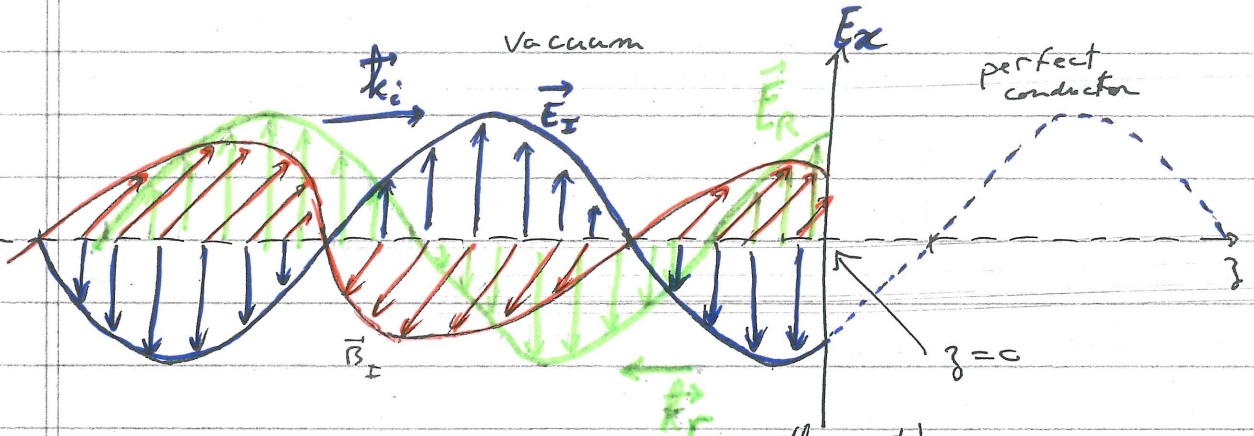


Standing waves:

Consider a plane wave incident on a perfect conducting mirror (normal incidence):

Standing waves

Consider a plane wave incident a perfect conducting mirror (normal incidence)



Incident plane wave:

$$\vec{E}_I = E_0 e^{i(kz - \omega t)} \hat{x}$$

$$\vec{B}_I = B_0 e^{i(kz - \omega t)} \hat{y}$$

Reflected plane wave:

$$\vec{E}_R = -E_0 e^{i(-kz - \omega t)} \hat{x}$$

$$\vec{B}_R = B_0 e^{i(-kz - \omega t)} \hat{y}$$

Stopped here ($r = \frac{E_0}{E_0} = -1$)

$$\Rightarrow \vec{E}_{\text{total}} = \vec{E}_I + \vec{E}_R = E_0 \left[e^{i(kz - \omega t)} - e^{i(-kz - \omega t)} \right] \hat{x}$$

$$= E_0 e^{-i\omega t} \left[\frac{e^{ikz} - e^{-ikz}}{2i} \right] \hat{x}$$

$$= E_0 2i e^{-i\omega t} \sin(kz) \hat{x}$$

take the real part $\Rightarrow \vec{E}_{\text{total}} = 2E_0 \sin(kz) \sin(\omega t) \hat{x}$

$$\begin{aligned}
 \text{Also, } \vec{B}_{\text{total}} &= \vec{B}_I + \vec{B}_R = B_0 \left[e^{i(kz - \omega t)} + e^{i(-kz - \omega t)} \right] \hat{y} \\
 &= B_0 e^{-i\omega t} 2 \left[\frac{e^{ikz} + e^{-ikz}}{2} \right] \hat{y} \\
 &= 2 B_0 e^{-i\omega t} \cos(kz) \hat{y}
 \end{aligned}$$

→ take real part ⇒ $\boxed{|\vec{B}_{\text{total}}| = 2 B_0 \cos(kz) \cos(\omega t) \hat{y}}$

- Results:
- 1) \vec{E} & \vec{B} fields are $\frac{\pi}{2}$ out of phase in space and time.
 - 2) \vec{E} & \vec{B} fields are standing waves (they don't travel).

