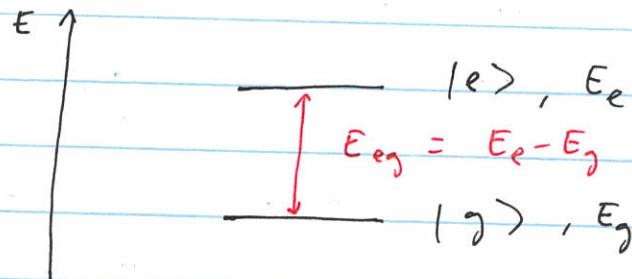


Tuesday, February 20, 2024

2-level atoms

(I) Definition: A 2-level atom is a quantum mechanical system with 2 distinct eigenstates & eigenenergies.



The basic Hamiltonian for the system is

$$H_0 = \begin{pmatrix} |g\rangle & |e\rangle \\ \langle g| & \langle e| \end{pmatrix} \begin{pmatrix} E_g & 0 \\ 0 & E_e \end{pmatrix}$$

One can write an arbitrary state of the system as

$$|\psi\rangle = \alpha|g\rangle + \beta|e\rangle \equiv (\alpha|g\rangle + \beta|e\rangle)^{\frac{1}{2}}$$

normalization requires

$$|\langle\psi|\psi\rangle|^2 = |\alpha|^2 + |\beta|^2 = 1$$

(II)

Bloch Sphere picture

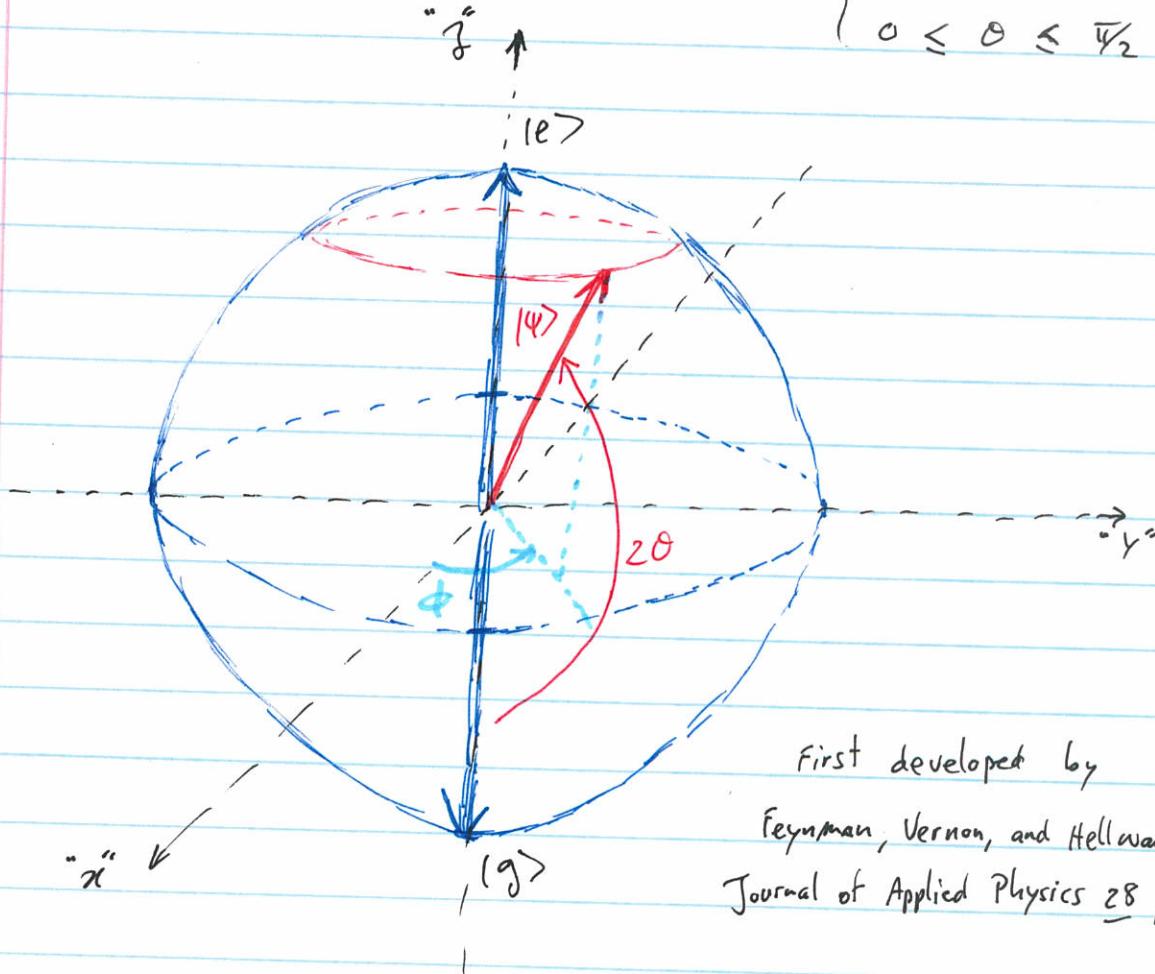
We want a geometric representation of the state of a 2-level system.

→ we could use a 2D vector space representation
 ↳ but phase is hard to show then!

A 2-level system is like a spin- $\frac{1}{2}$ system with
 $| \uparrow \rangle = | e \rangle$ and $| \downarrow \rangle = | g \rangle$

we note that we can write $| \psi \rangle = \cos\theta | g \rangle + e^{i\phi} \sin\theta | e \rangle$

$$\begin{cases} 0 \leq \phi < 2\pi \\ 0 \leq \theta \leq \frac{\pi}{2} \end{cases}$$



III Time-dependence

What happens when you let an arbitrary state evolve in time?

$$\text{Easy: } |\psi\rangle = \alpha|g\rangle + \beta|e\rangle \Rightarrow |\psi(t)\rangle = \alpha e^{-i\frac{E_g}{\hbar}t}|g\rangle + \beta e^{-i\frac{E_e}{\hbar}t}|e\rangle$$

$$= e^{\underbrace{-i\frac{E_g}{\hbar}t}_{\text{overall phase}}} [\alpha|g\rangle + \beta e^{\underbrace{-i\frac{(E_e - E_g)}{\hbar}t}_{\text{relative phase}}}|e\rangle]$$

\hookrightarrow not important \hookrightarrow totally important

$$P(|g\rangle)_{|\psi(t)\rangle} = |\langle g | \psi(t) \rangle|^2 = |\alpha|^2 \quad \left. \begin{array}{l} \text{no measurable} \\ \text{time-dependence} \end{array} \right\}$$

$$P(|e\rangle)_{|\psi(t)\rangle} = |\langle e | \psi(t) \rangle|^2 = |\beta|^2$$

[show phase precession on Bloch sphere]

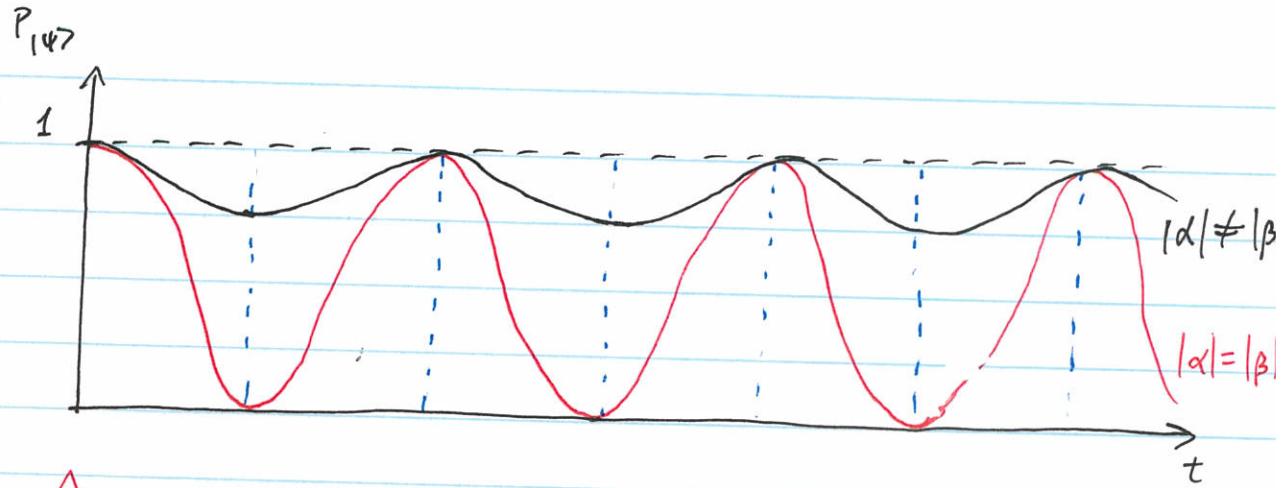
more interesting

$$P(|\psi\rangle)_{|\psi(t)\rangle} = |\langle \psi | \psi(t) \rangle|^2$$

$$= \left| (\alpha^* \langle g | + \beta^* \langle e |) (\alpha e^{-i\frac{E_g}{\hbar}t} |g\rangle + \beta e^{-i\frac{E_e}{\hbar}t} |e\rangle) \right|^2$$

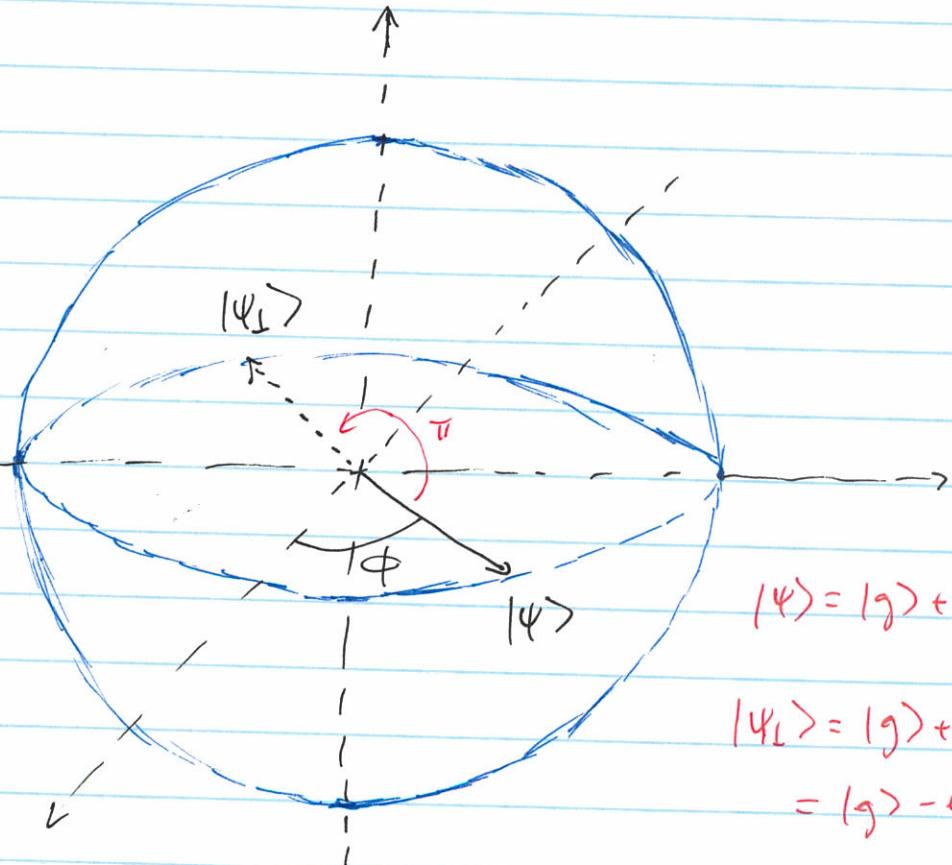
↓
(some algebra)

$$= \frac{1}{2} \left[1 + (|\alpha|^2 - |\beta|^2)^2 \right] + \frac{1}{2} \left[1 - (|\alpha|^2 - |\beta|^2)^2 \right] \cos \left(\frac{E_{eg}}{\hbar} t \right)$$



So there is time dependence if you look in another basis.

Bloch Sphere picture



$$\begin{aligned} |ψ_L\rangle &= |g\rangle + e^{i(\phi-\pi)} |e\rangle \\ &= |g\rangle - e^{i\phi} |e\rangle \end{aligned}$$

if $|\psi\rangle = \alpha|g\rangle + \beta|e\rangle$,
then $|\psi_L\rangle = \beta^*|g\rangle - \alpha^*|e\rangle$

2-level atom in an EM field

QM

Classical EM

(I) Hamiltonian

The interaction between an electric field and the electric dipole moment of an atom is given by

$$H_{\text{int}} = -\vec{d} \cdot \vec{E} = e \vec{R} \cdot \vec{E}_0 \cos(\omega_e t) \quad \text{with } \vec{z} = 0$$

(see Griffiths
eqn. 4.6)

(Assumption: single e^- orbiting a nucleus)

In the 2-level atom basis $\{|g\rangle, |e\rangle\}$, we have

$$H_{\text{int}} = e \vec{E}_0 \cos(\omega_e t) \begin{bmatrix} \langle g | \vec{R} | g \rangle & \langle g | \vec{R} | e \rangle \\ \langle e | \vec{R} | g \rangle & \langle e | \vec{R} | e \rangle \end{bmatrix}$$

↑
parity operator for

Recall: $[P, H_{\text{atom}}] = 0 \Rightarrow |g\rangle \text{ and } |e\rangle \text{ have a definite parity.}$

$$\langle g | \vec{R} | g \rangle = \sum_{\text{even}} \langle g | (x, y, z) | g \rangle = \int \psi_g^*(x, y, z) \psi_g d^3r$$

similarly, $\langle e | \vec{R} | e \rangle = 0$

odd odd odd

= 0

"2-level atom" atom = alkali atom $nS \leftrightarrow nP$ transition
 (1 valence e^-) $|g\rangle \leftrightarrow |e\rangle$

$$\begin{aligned} \text{In AMO, we define } \mathcal{R} &= \frac{e \langle g | \vec{r} | e \rangle}{\hbar} \cdot \vec{E}_0 \\ &= \frac{\langle g | H_{\text{int}} | e \rangle}{\hbar} \\ &= \text{Rabi frequency} \end{aligned}$$

So we have:

$$H = \underbrace{\begin{pmatrix} E_g & 0 \\ 0 & E_e \end{pmatrix}}_{H_0} + \underbrace{\hbar \begin{pmatrix} 0 & \mathcal{R} \\ \mathcal{R}^* & 0 \end{pmatrix} \cos(\omega_r t)}_{H_{\text{int}}}$$