

Thursday, February 29, 2024

#1

Dressed atom Hamiltonian (includes off resonance case)

$$H = H_{\text{atom}} + H_{\text{photons}} + H_{\text{interaction}}$$

$$= \begin{pmatrix} E_g & 0 \\ 0 & E_c \end{pmatrix} + \hbar \omega_e \begin{pmatrix} N+1 & 0 \\ 0 & N \end{pmatrix} + \hbar \begin{pmatrix} 0 & \mathcal{R}/2 \\ \mathcal{R}^*/2 & 0 \end{pmatrix}$$

in the $\{|1\rangle, |2\rangle\}$ basis, i.e. $\{|g\rangle_a |N+1\rangle_{\omega_e}, |e\rangle_a |N\rangle_{\omega_e}\}$

$$H = \hbar \begin{pmatrix} \omega_{eg} + \delta & \mathcal{R}/2 \\ \mathcal{R}^*/2 & \omega_{eg} \end{pmatrix}$$

subtract " $\hbar \omega_e N$ " & E_g
from diagonal energies
($\delta = \omega_e - \omega_{eg}$)
[it's just an energy offset]

$$\Rightarrow H = \hbar \begin{pmatrix} \delta & \mathcal{R}/2 \\ \mathcal{R}^*/2 & 0 \end{pmatrix}$$

subtract " $\hbar \omega_{eg}$ " from
diagonal terms.

Note: A rigorous derivation requires going into the rotating frame
of the atom (not an approximation).

AC Stark shift

Let's solve for the eigenenergies of the Hamiltonian

$$\det\{H - \lambda I\} = 0 \Leftrightarrow \begin{vmatrix} \delta - \lambda & \omega_0 \\ \omega_0^* & 0 - \lambda \end{vmatrix} = 0$$

$$\Leftrightarrow (\delta - \lambda)(-\lambda) - \frac{|\omega_0|^2}{4} = 0$$

$$\Leftrightarrow \lambda^2 - \delta\lambda - \frac{|\omega_0|^2}{4} = 0$$

$$\Rightarrow \lambda_{\pm} = \frac{\delta \pm \sqrt{\delta^2 + |\omega_0|^2}}{2} = \frac{E_{\pm}}{\hbar}$$

for $\delta = 0$: $E_{\pm} = \pm \frac{\hbar |\omega_0|}{2}$ "on-resonance"

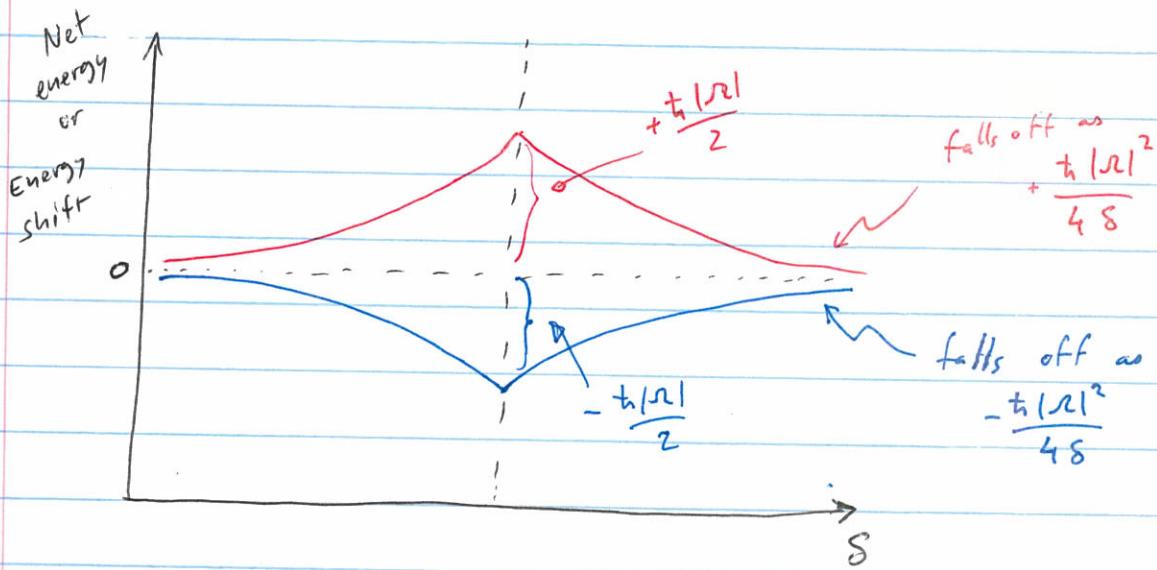
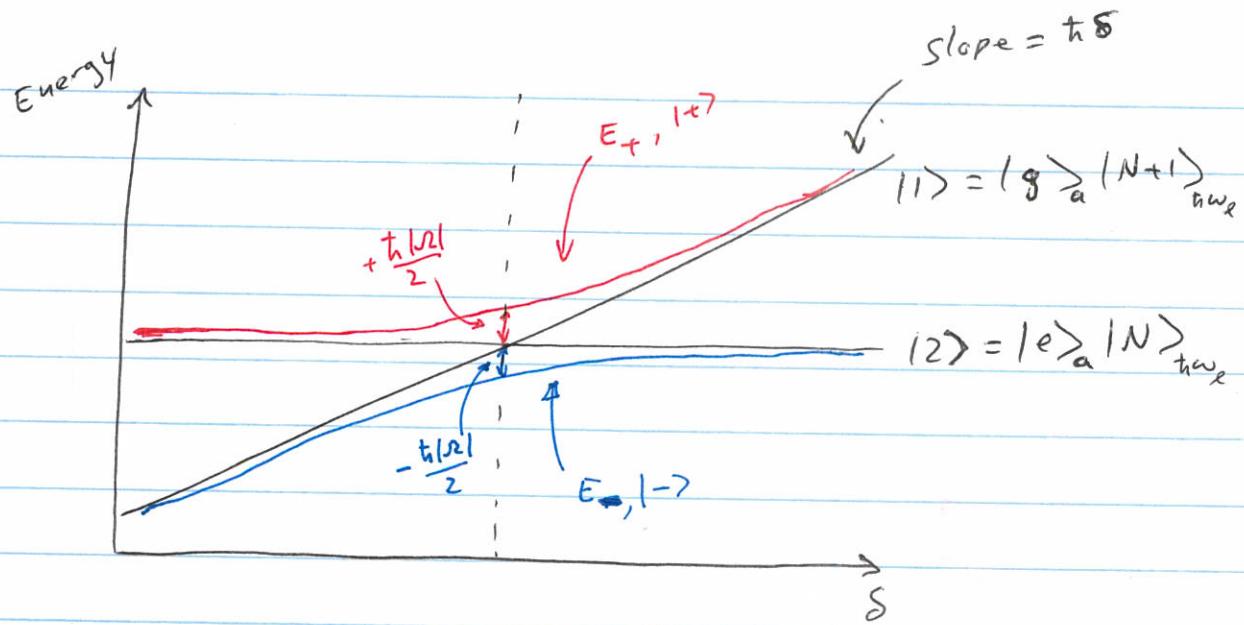
for $|\delta| \gg |\omega_0|$ "far off resonance"

$$\sqrt{\delta^2 + |\omega_0|^2} \approx |\delta| + \frac{1}{2} \frac{|\omega_0|^2}{|\delta|}$$

Thus $E_{\pm} = \hbar \frac{\delta \pm \left(|\delta| + \frac{1}{2} \frac{|\omega_0|^2}{|\delta|} \right)}{2}$

$$\Rightarrow \begin{cases} E_+ = \frac{\hbar(\delta + |\delta|)}{2} + \frac{\hbar|\omega_0|^2}{4|\delta|} \\ E_- = \frac{\hbar(\delta - |\delta|)}{2} - \frac{\hbar|\omega_0|^2}{4|\delta|} \end{cases}$$

laser energy,
does not depend
on " ω_0 ".



Next, let's find the eigenstates (or "dressed states")

$$\frac{E_+}{t} = \lambda_+ = \frac{s + \alpha'_l}{2} \quad \text{where } \alpha'_l = \sqrt{s^2 + |\alpha_l|^2}$$

$$\begin{bmatrix} s - \lambda_+ & \alpha'_l/2 \\ \alpha'^*_l & 0 - \lambda_+ \end{bmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$\begin{pmatrix} \frac{s}{2} - \frac{s'z}{2} & s_2 \\ \frac{s}{2} & -\frac{s}{2} - \frac{s'z}{2} \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} |1\rangle \\ |2\rangle \end{pmatrix}$$

$$\Rightarrow |+\rangle = \frac{s}{2} |1\rangle - \left(\frac{s}{2} - \frac{s'z}{2} \right) |2\rangle$$

with normalization, we can write $|+\rangle = \cos\theta |1\rangle + \sin\theta |2\rangle$
(for $s, s' \in \mathbb{R}$)

$$\text{similarly, } |- \rangle = -\sin\theta |1\rangle + \cos\theta |2\rangle$$

$$\text{with } \cos\theta = \frac{s}{\sqrt{s^2 + (s'-s)^2}} \quad \text{and } \sin\theta = \frac{s' - s}{\sqrt{s^2 + (s'-s)^2}}$$

$$\text{note: } \cos 2\theta = \cos^2\theta - \sin^2\theta = -\frac{s}{s'}$$

Example: On resonance: $s=0, s'=s$

$$\cos\theta = \frac{1}{\sqrt{2}}, \sin\theta = \frac{1}{\sqrt{2}} \Rightarrow \cos 2\theta = 0 \\ \Rightarrow \theta = \pi/4$$

$$\Rightarrow \begin{cases} |+\rangle = \frac{1}{\sqrt{2}} |1\rangle + \frac{1}{\sqrt{2}} |2\rangle \\ |- \rangle = -\frac{1}{\sqrt{2}} |1\rangle + \frac{1}{\sqrt{2}} |2\rangle \end{cases}$$

These are the eigenstates of the atom-photon system, but we perform measurements in the $\{|1\rangle, |2\rangle\}$ basis typically, so we observe oscillations at $\omega = \frac{\Delta E}{\hbar} = \Omega$, i.e. Rabi flopping.

The atom-light interaction mixes the $|1\rangle$ and $|2\rangle$ states and lifts their degeneracy.

For EM field off : $|+\rangle = |1\rangle$
 $|-\rangle = |2\rangle$