

Tuesday, March 26, 2024

Resonance Fluorescence - Density matrix approach

coherent evolution:  $i\hbar \frac{d}{dt} \rho = [H, \rho]$

↑ density matrix

with  $H = \begin{pmatrix} E_g & \hbar\Omega \cos(\omega_L t) \\ \hbar\Omega^* \cos(\omega_L t) & E_e \end{pmatrix} + RWA$

or  $H_{\text{dressed}} = \begin{pmatrix} S & \Omega/2 \\ \Omega^*/2 & 0 \end{pmatrix}$

Add in decoherence / spontaneous emission (decay rate  $\gamma$ ):

$$\left\{ \begin{aligned} \frac{d}{dt} \rho_{gg} &= -\frac{i}{\hbar} [H, \rho]_{gg} + \gamma \rho_{ee} \\ \frac{d}{dt} \rho_{ee} &= -\frac{i}{\hbar} [H, \rho]_{ee} - \gamma \rho_{ee} \\ \frac{d}{dt} \rho_{eg} &= -\frac{i}{\hbar} [H, \rho]_{eg} - \frac{1}{2} \gamma \rho_{eg} \\ \frac{d}{dt} \rho_{ge} &= -\frac{i}{\hbar} [H, \rho]_{ge} - \frac{1}{2} \gamma \rho_{ge} \end{aligned} \right.$$

optical Bloch equations

⇒ Solve for

$$\begin{matrix} \rho_{gg}, \rho_{ee}, \rho_{eg}, \rho_{ge} \\ \uparrow \quad \uparrow \\ P_g \quad P_e \end{matrix}$$

all continuous quantities (wavefunction collapse is already included)

Marquardt, Robbins & Hellberg

For generalization to multi-level atoms and multiple lasers see JOSA B, 13, 1384 (1996)

We are interested in the steady state scattering rate:

$$\text{set } \frac{d}{dt} \rho = 0 = -\frac{i}{\hbar} [H, \rho]_{ij} \quad \text{+ } \delta f_{ij} \quad \left(\frac{1}{2}\right)$$

choose appropriately

scattering rate =  $\delta_{\text{scattering}} = \delta P_e = \delta f_{ee}$

↑ probability to be in the excited state.

$$P_g = |c_g|^2 = f_{gg}$$

$$P_e = |c_e|^2 = f_{ee}$$

some work ↓

$$= \frac{\rho_0}{1 + \rho_0 + \left(\frac{2S}{\delta}\right)^2} \frac{\delta}{2} = \delta_{\text{scattering}}$$

where  $\rho_0 = \text{saturation parameter} = \frac{I}{I_{\text{sat}}} = \frac{2 |R|^2}{\gamma^2}$

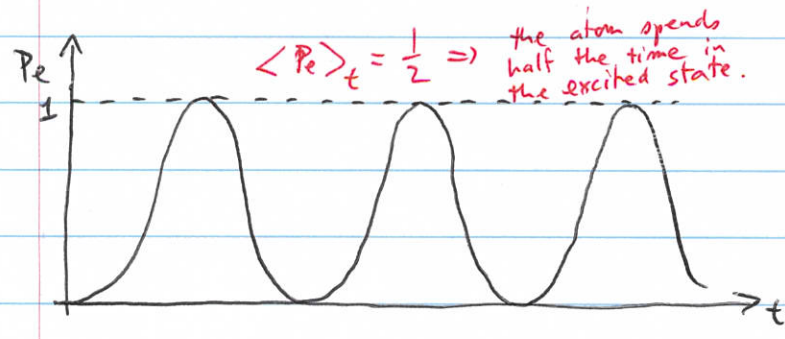
$I_{\text{sat}} = \frac{\pi \hbar c}{3 \lambda^2 \tau} = 1.6 \text{ mW/cm}^2$  for  $^{87}\text{Rb}$  ( $D_2$  line)

= intensity required to reach half of maximum upper state population in steady state (i.e.  $\frac{1}{4}$ )

On resonance:  $\delta = 0$

Case  $I \gg I_{\text{sat}}$ , i.e.  $\rho_0 \rightarrow +\infty$ , then  $\delta_{\text{scattering}} = \frac{\delta}{2} = 19 \times 10^6 \frac{\text{photons}}{\text{s}}$

rubidium



In strong contradiction with classical physics, but in good agreement with experiment

Case  $I \ll I_{\text{sat}}$ , i.e.  $\rho_0 \rightarrow 0$ , then  $\gamma_{\text{scattering}} \approx \rho_0 \frac{\gamma}{2} \propto \text{Incident Intensity}$

good agreement with classical physics

off-Resonance

Linewidth for  $\rho_0 \rightarrow 0$ :  $\gamma_s \approx \frac{1}{1 + \left(\frac{2\delta}{\gamma}\right)^2} \rho_0 \frac{\gamma}{2}$

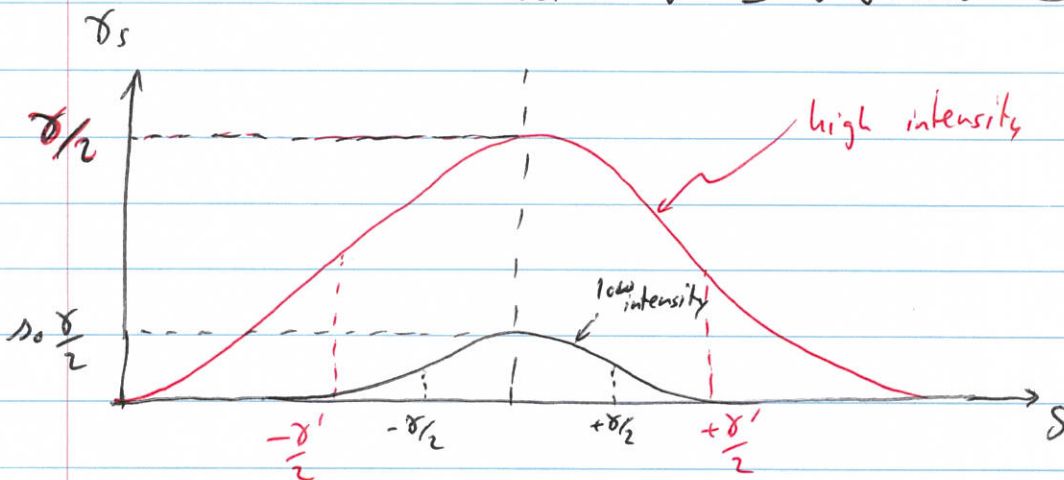
$\Rightarrow \text{FWHM} = \gamma$

Linewidth for  $\rho_0 \rightarrow +\infty$

"power broadening"

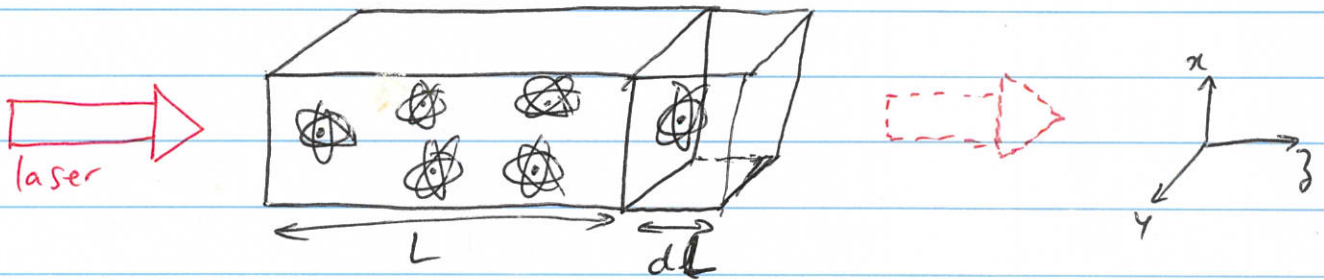
$$\gamma_s = \frac{\rho_0}{1 + \rho_0} \frac{1}{1 + \left(\frac{2\delta}{\gamma\sqrt{1+\rho_0}}\right)^2} \frac{\gamma}{2} \approx \frac{1}{1 + \left(\frac{2\delta}{\gamma'}\right)^2} \frac{\gamma}{2}$$

with  $\gamma' = \gamma\sqrt{1+\rho_0} = \text{FWHM}$



## Beer's law : Absorption in an atomic vapor

How is a laser attenuated as it travels through a gas of resonant atoms?



$$\text{gas density} = n = \frac{N_{\text{atoms}}}{\text{Volume}}$$

$$\text{change in intensity} = dI = \frac{d\text{Power}}{\text{Area}} = - \frac{\hbar \omega_p \gamma_s}{A} n A dl$$

$A =$  area of laser

photons  
per atom

# atoms in  
"dl" volume

$$\text{thus } \frac{dI}{dL} = -\hbar \omega_p \gamma_s n$$

In the weak intensity limit:  $\gamma_s = \frac{I}{I_{\text{sat}}} \cdot \frac{\gamma}{2}$  (on resonance)

$$\frac{dI}{dL} = -\hbar \omega_p \frac{\gamma}{2} \frac{n}{I_{\text{sat}}} I = - \frac{3\lambda^2}{2\pi} n I \quad (\text{on resonance})$$

$$= -\sigma_{\text{eg}} n I$$

$\sigma_{\text{eg}} =$  cross section  
on resonance

$$\Rightarrow I(z) = I_0 e^{-\sigma_{eg} n z}$$

Beer's Law

note: optical depth =  $\sigma_{eg} n z$

$$\hookrightarrow I = I_0 e^{-\text{optical depth}}$$

Application  $\hat{=}$  absorption imaging of cold atoms.