

Monday, February 8, 2021

VI Spherical Coordinates (continued)

Volume element: $dV = r^2 \sin \theta dr d\theta d\phi$

↳ replace $dx dy dz$ with this expression.

⚠ note: $dx dy dz \neq r^2 \sin \theta dr d\theta d\phi$

more generally: $dV = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} & \frac{\partial x}{\partial \phi} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} & \frac{\partial y}{\partial \phi} \\ \frac{\partial z}{\partial r} & \frac{\partial z}{\partial \theta} & \frac{\partial z}{\partial \phi} \end{vmatrix} dr d\theta d\phi$

determinant of the Jacobian matrix = $r^2 \sin \theta$

for cylindrical coordinates: $dV = r dr dz d\phi$

Grad, Div, Curl, etc

Gradient: $\vec{\nabla} f(\vec{r}) = \frac{\partial f}{\partial x} \hat{x} + \frac{\partial f}{\partial y} \hat{y} + \frac{\partial f}{\partial z} \hat{z}$

lots of algebra

$= \frac{\partial f}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial f}{\partial \theta} \hat{\theta} + \frac{1}{r \sin \theta} \frac{\partial f}{\partial \phi} \hat{\phi}$

↑ depends on location!

$$\underline{\text{ex:}} \quad \frac{\partial f}{\partial x} = \frac{\partial f}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial f}{\partial \theta} \frac{\partial \theta}{\partial x} + \frac{\partial f}{\partial \phi} \frac{\partial \phi}{\partial x}$$

$$\frac{1}{2} \frac{2x}{\sqrt{x^2 + y^2 + z^2}} = \frac{x}{r} = \sin \theta \cos \phi$$

also ~~unit~~ $\hat{n} = \sin \theta \cos \phi \hat{r} + \cos \theta \cos \phi \hat{\theta} - \sin \phi \hat{\phi}$
 etc...

SPHERICAL AND CYLINDRICAL COORDINATES

#3

Spherical

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases} \quad \begin{cases} \hat{x} = \sin \theta \cos \phi \hat{r} + \cos \theta \cos \phi \hat{\theta} - \sin \phi \hat{\phi} \\ \hat{y} = \sin \theta \sin \phi \hat{r} + \cos \theta \sin \phi \hat{\theta} + \cos \phi \hat{\phi} \\ \hat{z} = \cos \theta \hat{r} - \sin \theta \hat{\theta} \end{cases}$$

$$\begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \theta = \tan^{-1} \left(\frac{\sqrt{x^2 + y^2}}{z} \right) \\ \phi = \tan^{-1} (y/x) \end{cases} \quad \begin{cases} \hat{r} = \sin \theta \cos \phi \hat{x} + \sin \theta \sin \phi \hat{y} + \cos \theta \hat{z} \\ \hat{\theta} = \cos \theta \cos \phi \hat{x} + \cos \theta \sin \phi \hat{y} - \sin \theta \hat{z} \\ \hat{\phi} = -\sin \phi \hat{x} + \cos \phi \hat{y} \end{cases}$$

Cylindrical

$$\begin{cases} x = s \cos \phi \\ y = s \sin \phi \\ z = z \end{cases} \quad \begin{cases} \hat{x} = \cos \phi \hat{s} - \sin \phi \hat{\phi} \\ \hat{y} = \sin \phi \hat{s} + \cos \phi \hat{\phi} \\ \hat{z} = \hat{z} \end{cases}$$

$$\begin{cases} s = \sqrt{x^2 + y^2} \\ \phi = \tan^{-1} (y/x) \\ z = z \end{cases} \quad \begin{cases} \hat{s} = \cos \phi \hat{x} + \sin \phi \hat{y} \\ \hat{\phi} = -\sin \phi \hat{x} + \cos \phi \hat{y} \\ \hat{z} = \hat{z} \end{cases}$$

Spherical. $d\mathbf{l} = dr \hat{r} + r d\theta \hat{\theta} + r \sin \theta d\phi \hat{\phi}; \quad d\tau = r^2 \sin \theta dr d\theta d\phi$

Gradient: $\nabla t = \frac{\partial t}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\theta} + \frac{1}{r \sin \theta} \frac{\partial t}{\partial \phi} \hat{\phi}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$

Curl: $\nabla \times \mathbf{v} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{r}$
 $+ \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\theta} + \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\phi}$

Laplacian: $\nabla^2 t = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial t}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial t}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 t}{\partial \phi^2}$

Cylindrical. $d\mathbf{l} = ds \hat{s} + s d\phi \hat{\phi} + dz \hat{z}; \quad d\tau = s ds d\phi dz$

Gradient: $\nabla t = \frac{\partial t}{\partial s} \hat{s} + \frac{1}{s} \frac{\partial t}{\partial \phi} \hat{\phi} + \frac{\partial t}{\partial z} \hat{z}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$

Curl: $\nabla \times \mathbf{v} = \left[\frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right] \hat{s} + \left[\frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right] \hat{\phi} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{z}$

Laplacian: $\nabla^2 t = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial t}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 t}{\partial \phi^2} + \frac{\partial^2 t}{\partial z^2}$

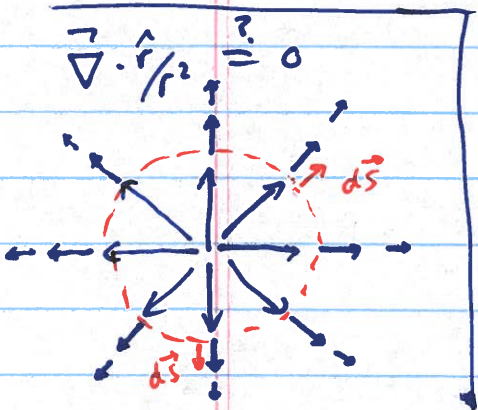
VII Divergence of $\vec{A} = \hat{r}/r^2$ [Δ plays a central role] in E & M

Divergence in spherical coordinates:

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$$

$$= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{1}{r^2} \right) = \frac{1}{r^2} \frac{\partial}{\partial r} (1) = \frac{0}{r^2}$$

= 0 ! Wrong



"Paradox": $\int_V (\vec{\nabla} \cdot \vec{A}) dV = \int_S \vec{A} \cdot d\vec{s} = 4\pi R^2 \frac{1}{r^2} \uparrow R$

$V =$ sphere of radius R $S =$ surface of sphere

$\int_V "0" dV = 0$

\neq

4π

does not depend on "R"

Resolution: $\vec{\nabla} \cdot \frac{\hat{r}}{r^2} = 4\pi \delta^3(\vec{r})$

even if $R = \epsilon$ "infinitesimal sphere" centered on origin

where $\delta^3(\vec{r}) = \delta(x) \delta(y) \delta(z)$

note: $\vec{\nabla} \cdot \left(\underbrace{\vec{\nabla} \frac{1}{r}}_{-\frac{\vec{r}}{r^2}} \right) = \nabla^2 \frac{1}{r} = -4\pi \delta^3(\vec{r})$

VIII Helmholtz Theorem

Short version: If you know $\vec{\nabla} \cdot \vec{F}$ and $\vec{\nabla} \times \vec{F}$, as well as the value of \vec{F} on some boundary (e.g. $\vec{F}(r \rightarrow +\infty) \rightarrow 0$) then \vec{F} is uniquely determined.

Long version:

An arbitrary vector field $\vec{F}(\vec{r})$ can always be decomposed into a sum of 2 vector fields such that

$$\vec{F}(\vec{r}) = \vec{F}_\perp + \vec{F}_\parallel$$

with $\vec{\nabla} \cdot \vec{F}_\perp = 0$ and $\vec{\nabla} \times \vec{F}_\parallel = \vec{0}$

and more precisely as curl + gradient

$$\vec{F}(\vec{r}) = \underbrace{\vec{\nabla} \times \vec{U}(\vec{r})}_{\vec{F}_\perp} - \underbrace{\vec{\nabla} \Omega(\vec{r})}_{\vec{F}_\parallel}$$

So long as the integrals converge, then \vec{U} and Ω are "uniquely" given by

$$\Omega(\vec{r}) = \frac{1}{4\pi} \int d^3r' \frac{\vec{\nabla}_{\vec{r}'} \cdot \vec{F}(\vec{r}')}{|\vec{r} - \vec{r}'|} = \text{scalar potential}$$

$$\vec{U}(\vec{r}) = \frac{1}{4\pi} \int d^3r' \frac{\vec{\nabla}_{\vec{r}'} \times \vec{F}(\vec{r}')}{|\vec{r} - \vec{r}'|} = \text{vector potential}$$

note: $\vec{F}(\vec{r})$ must fall off faster than $1/r$ for $r \rightarrow +\infty$

proof: see PHYS 610 lecture notes (Spring 2018) for March 1, 2018

$$\hookrightarrow \text{proof uses } \nabla^2 \frac{1}{r} = -4\pi \delta^3(\vec{r})$$

Corollary 1: If $\vec{\nabla} \times \vec{F} = 0$, then $\vec{F}(\vec{r})$ can be derived from a ^{scalar} potential \rightarrow electrostatic potential V for ~~E~~ electric field \vec{E} .

Corollary 2: If $\vec{\nabla} \cdot \vec{F} = 0$, then $\vec{F}(\vec{r})$ can be derived from a vector potential. \rightarrow magnetic vector potential \vec{A} for the magnetic field \vec{B} .