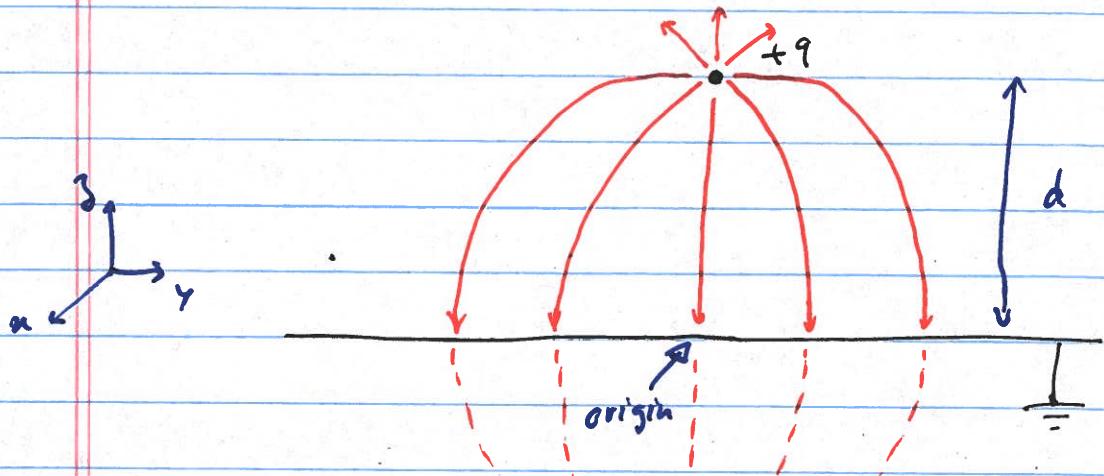


Wednesday, March 3, 2021

## METHOD OF IMAGES (continued)

point charge above an infinite conducting grounded plane.



$$\text{Solution: } V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \left[ \frac{q}{\sqrt{x^2 + y^2 + (z-d)^2}} - \frac{q}{\sqrt{x^2 + y^2 + (z+d)^2}} \right]$$

$\hookrightarrow$  satisfies b.c.  $V(z=0) = 0$  } the solution is only valid  
 $V(r \rightarrow \infty) = 0$  } in the region without the  
image charge.

Q: What's the surface charge density?

$$\Delta E_{\text{surface}} = \vec{E}_{\text{ext}} - \underbrace{\vec{E}_c}_{=0} = \frac{\sigma}{\epsilon_0} \hat{n} \Leftrightarrow \vec{E}(z=0^+) = \frac{\sigma}{\epsilon_0} \hat{n}$$

$$\Rightarrow -\nabla V \Big|_{z=0^+} = \frac{\sigma}{\epsilon_0} \hat{n} \Leftrightarrow -\frac{\partial V}{\partial z} \Big|_{z=0^+} \hat{z} = \frac{\sigma}{\epsilon_0} \hat{z} \quad \left[ \frac{\partial V}{\partial x} \Big|_{z=0^+} = \frac{\partial V}{\partial y} \Big|_{z=0^+} = 0 \right]$$

$$\left. \frac{\partial V}{\partial x} \right|_{y=0} = \frac{1}{4\pi\epsilon_0} q \left[ -\frac{1}{2} \frac{2x}{[x^2+y^2+(z-d)^2]^{3/2}} - \left( -\frac{1}{2} \right) \frac{2x}{[x^2+y^2+(z+d)^2]^{3/2}} \right]$$

$$= 0 \quad \text{also} \quad \left. \frac{\partial V}{\partial y} \right|_{y=0} = 0$$

$$\left. \frac{\partial V}{\partial z} \right|_{y=0} = \frac{1}{4\pi\epsilon_0} q \left[ \left( -\frac{1}{2} \right) \frac{2(z-d)}{[x^2+y^2+(z-d)^2]^{3/2}} - \left( -\frac{1}{2} \right) \frac{2(z+d)}{[x^2+y^2+(z+d)^2]^{3/2}} \right]$$

$$= \frac{1}{4\pi\epsilon_0} \frac{2d}{[x^2+y^2+d^2]^{3/2}}$$

since  $\left. \frac{\partial V}{\partial z} \right|_{y=0} = -\frac{\sigma}{\epsilon} \Rightarrow \sigma = -\frac{1}{4\pi\epsilon_0} q \frac{(2d)}{[x^2+y^2+d^2]^{3/2}}$

$$\Rightarrow \boxed{\sigma = -\frac{q}{2\pi} \frac{d}{[x^2+y^2+d^2]^{3/2}}}$$

induced charge density

total induced charge =  $q_c = \int \sigma d\sigma$

$$= -\frac{q}{2\pi} d \int \frac{dz dy}{[x^2+y^2+d^2]^{3/2}} = -\frac{q d}{2\pi} \int_0^\infty \int_0^{2\pi} \frac{r dr d\phi}{[r^2+d^2]^{3/2}}$$

$x, y$   
plane

$$= -\frac{q}{2\pi} d 2\pi \int_0^\infty \frac{r dr}{[r^2+d^2]^{3/2}} = -q d \left( \frac{-1}{\sqrt{r^2+d^2}} \right)_0^\infty$$

$$= -q d \left( 0 - \left( \frac{1}{d} \right) \right) = -q$$

Thus the induced charge =  $q_c = -q$  (as expected)

Force on the charge  $q$  by the conducting plane

$$F = \frac{1}{4\pi\epsilon_0} \frac{(-q)}{(2d)^2} q \hat{j}$$

Coulomb force

due to image  $\Rightarrow$  charge is attracted  
charge to the conducting  
plane.

Electrostatic energy, is half that for the charge + image charge  
half of universe is missing

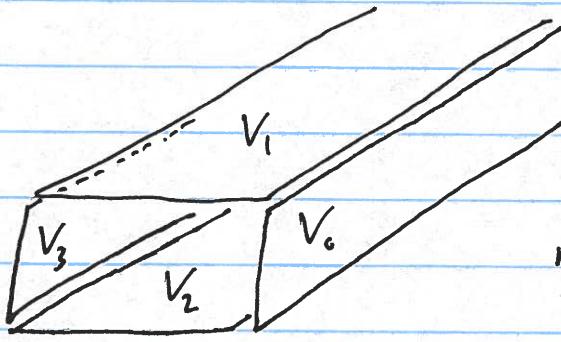
$$\begin{aligned} W &= \int_{-\infty}^d \vec{F} \cdot d\vec{l} = \frac{1}{4\pi\epsilon_0} \int_{-\infty}^d \frac{q^2}{4j^2} dj \\ &= \frac{q^2}{4\pi\epsilon_0} \left[ \frac{1}{4} j - \frac{1}{3} j^3 \right]_{-\infty}^d = \frac{q^2}{4\pi\epsilon_0} \frac{1}{4} \left( -\frac{1}{d} - 0 \right) \\ &= -\frac{1}{4\pi\epsilon_0} \frac{q^2}{4d} \end{aligned}$$

$$\Rightarrow \boxed{W = -\frac{1}{4\pi\epsilon_0} \frac{q^2}{4d}}$$

## Separation of Variables

Consider a system with Cartesian symmetry  
(in the geometry)

Example:



rectangular duct with  
conducting walls  
(no extra charges)

Q: What is the potential inside the duct?

no charges (except on boundaries)

→ Laplace's equation:  $\nabla^2 V = 0$

$$\Rightarrow \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) V(x, y, z) = 0$$

Ausatz: We will consider separable solutions of the form

only possible because of  
symmetry  
(each coordinate gets  
its own function)

$$V(x, y, z) = X(x)Y(y)Z(z)$$

$$\Rightarrow Y(y)Z(z) \frac{\partial^2}{\partial x^2} X + X(x)Z(z) \frac{\partial^2}{\partial y^2} Y + X(x)Y(y) \frac{\partial^2}{\partial z^2} Z = 0$$

Divide by  $X(x)Y(y)Z(z)$

$$\Rightarrow \frac{1}{X(x)} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y(y)} \frac{\partial^2 Y}{\partial y^2} + \frac{1}{Z(z)} \frac{\partial^2 Z}{\partial z^2} = 0$$



Since the solution must work for all  $(x, y, z)$  in the solution volume, then each term is equal to a constant such that

$$\Rightarrow \left\{ \begin{array}{l} \frac{1}{X(x)} \frac{\partial^2 X}{\partial x^2} = C_x \Leftrightarrow \frac{\partial^2 X}{\partial x^2} = C_x X = \alpha^2 X \\ \frac{1}{Y(y)} \frac{\partial^2 Y}{\partial y^2} = C_y \Leftrightarrow \frac{\partial^2 Y}{\partial y^2} = C_y Y = \beta^2 Y \\ \frac{1}{Z(z)} \frac{\partial^2 Z}{\partial z^2} = C_z \Leftrightarrow \frac{\partial^2 Z}{\partial z^2} = C_z Z = \gamma^2 Z \end{array} \right.$$

with  $C_x + C_y + C_z = 0 \Leftrightarrow \alpha^2 + \beta^2 + \gamma^2 = 0$

$$\alpha, \beta, \gamma \in \mathbb{C}$$

Solutions:  $X_\alpha = \begin{cases} A_0 + B_0 x & \alpha = 0 \\ A_\alpha e^{\alpha x} + B_\alpha e^{-\alpha x} & \alpha \neq 0 \end{cases}$

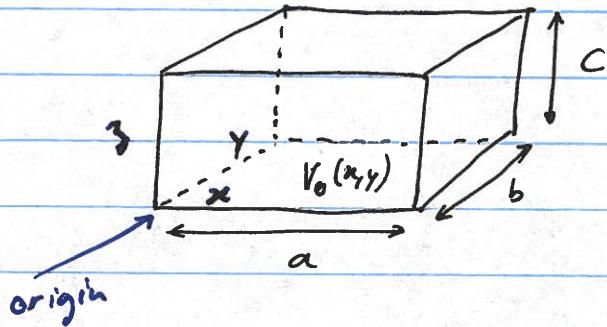
$$Y_\beta = \begin{cases} C_0 + D_0 y & \beta = 0 \\ C_\beta e^{\beta y} + D_\beta e^{-\beta y} & \beta \neq 0 \end{cases}$$

$$Z_\gamma = \begin{cases} E_0 + F_0 z & \gamma = 0 \\ E_\gamma e^{\gamma z} + F_\gamma e^{-\gamma z} & \gamma \neq 0 \end{cases}$$

General Solution: We consider any linear superposition of these solutions.

$$V(x, y, z) = \sum_{\alpha, \beta, \gamma} X_\alpha(x) Y_\beta(y) Z_\gamma(z) \underbrace{\delta(\alpha^2 + \beta^2 + \gamma^2)}_{\text{enforces } \alpha^2 + \beta^2 + \gamma^2 = 0}$$

Example: Consider a rectangular conducting box with 5 sides at potential  $V=0$  and one side at potential  $V_0(x, y)$ .



## boundary conditions:

$$\left\{ \begin{array}{l} X_\alpha(x=0) = 0 \\ X_\alpha(x=a) = 0 \end{array} \right. \quad \left\{ \begin{array}{l} Y_\beta(y=0) = 0 \\ Y_\beta(y=b) = 0 \end{array} \right.$$

$$Z_\gamma (\beta = c) = 0 \quad \text{and} \quad V(x, y, 0) = V_0(x, y)$$

Easier      Harder

For  $X_\alpha$ :  $X_\alpha(n=0) = 0 \Rightarrow A_\alpha + B_\alpha = 0 \Leftrightarrow A_\alpha = -B_\alpha$

$$X_\alpha(\kappa=a) = 0 \Rightarrow A_\alpha e^{\alpha a} - A_\alpha e^{-\alpha a} = 0$$

$$\Leftrightarrow e^{\alpha a} = e^{-\beta a}$$

$$\Leftrightarrow e^{2\alpha a} = 1 \Rightarrow 2\alpha a = 0 \Rightarrow \alpha = 0$$

or

$$2\alpha a = n \cdot 2\pi$$

$$\Rightarrow \lambda = \frac{i\pi}{a}$$

with  $n = 0, \pm 1, \pm 2, \dots$