

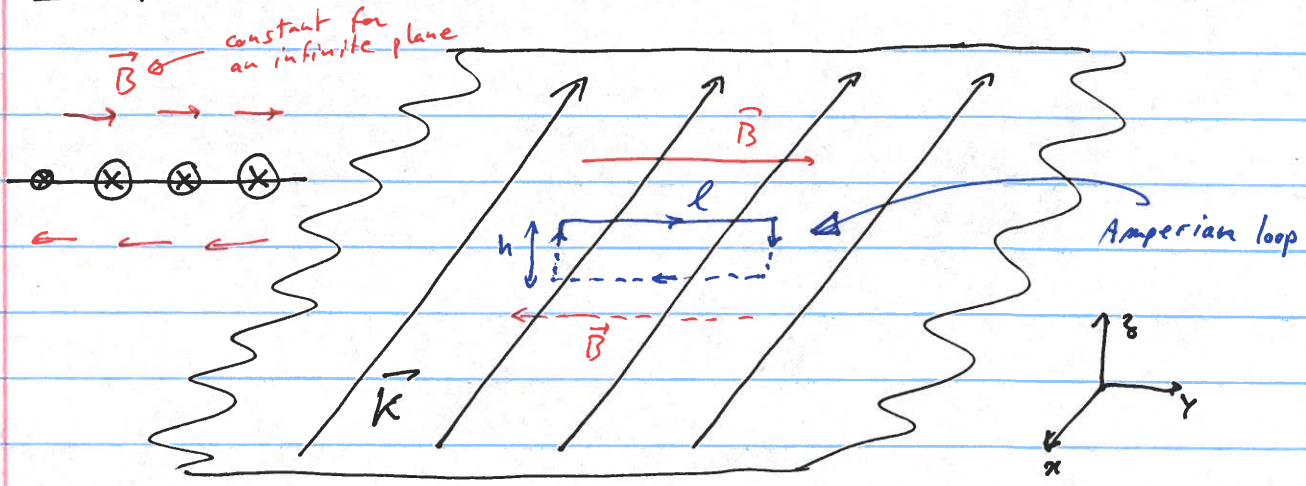
Monday, April 17, 2023

Ampère's Law (continued) [chpt 5.3.3]

$$\nabla \times \vec{B}(\vec{r}) = \mu_0 \vec{J}(\vec{r}) \Leftrightarrow \oint_{\text{loop}} \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enclosed}}$$

↑ magnetic field ↑ current density

Example: Planar Surface Current



$$\oint_{\text{loop}} \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enclosed}}$$

$$\Leftrightarrow \underbrace{\vec{B} \cdot \vec{l}}_{\parallel} + \underbrace{\vec{B} \cdot \vec{h}}_{=0} + \underbrace{\vec{B} \cdot \vec{l}}_{\parallel} + \underbrace{\vec{B} \cdot \vec{h}}_{=0} = \mu_0 k l$$

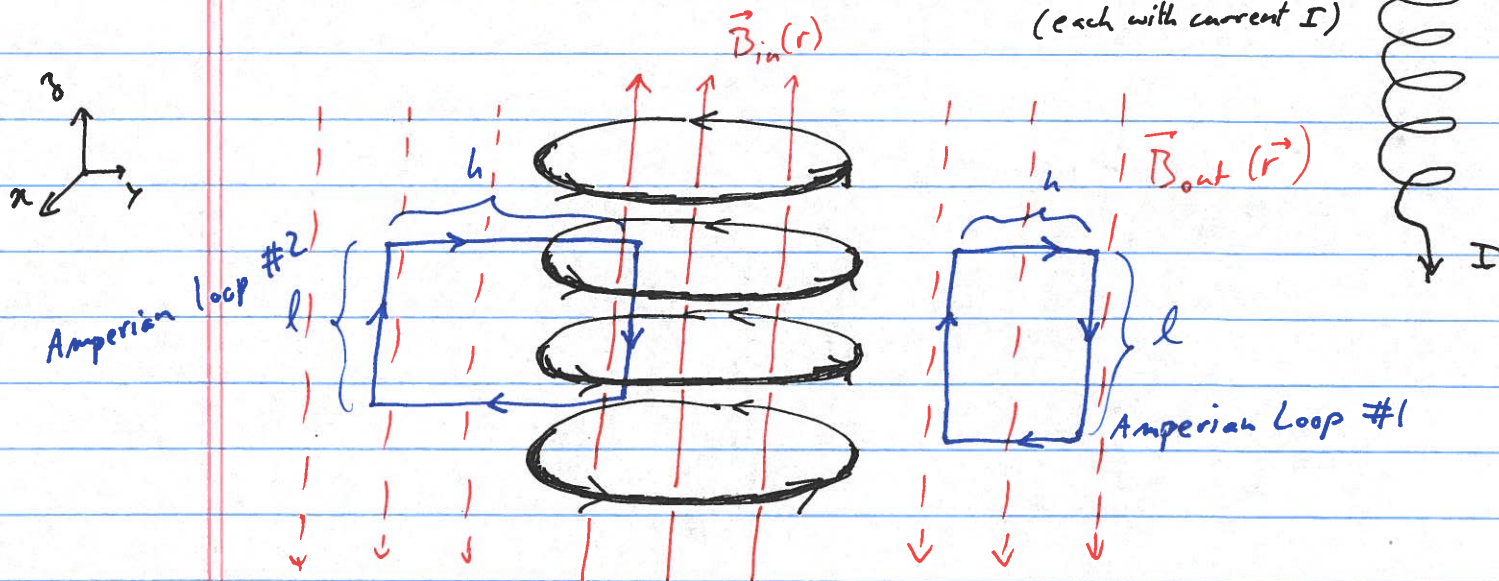
$\vec{B} \perp \vec{h}$

$$\Leftrightarrow 2Bl = \mu_0 k l \Leftrightarrow \begin{cases} \vec{B} = \frac{\mu_0 k}{2} \hat{y} & \text{above} \\ \vec{B} = -\frac{\mu_0 k}{2} \hat{y} & \text{below} \end{cases}$$

note: B-field is independent of height.
(for an infinite plane)

Example: Ideal Solenoid ("infinitely" long)

→ Approximate solenoid as a series of independent current loops. (each with current I)



Amperian loop #1.

$$\oint_{\text{loop 1}} \vec{B} \cdot d\vec{\ell} = \mu_0 \underbrace{I_{\text{enclosed}}}_{=0}$$

$$\Leftrightarrow -\vec{B}_{\text{near}} \cdot \vec{\ell} + \underbrace{\int \vec{B} \cdot d\vec{\ell}}_{=0} + \vec{B}_{\text{far}} \cdot \vec{\ell} + \underbrace{\int \vec{B} \cdot d\vec{\ell}}_{=0} = 0$$

$$\Rightarrow \vec{B}_{\text{near}} = \vec{B}_{\text{far}} \quad (\Rightarrow) \quad \vec{B}_{\text{out}}(r) = \text{constant}$$

Very far from solenoid, we expect/require: $\vec{B}_{\text{out}}(r \rightarrow \infty) = 0$

$$\Rightarrow \boxed{\vec{B}_{\text{out}}(r) = 0}$$

Amperian Loop #2

$$\oint_{\text{Loop 2}} \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enclosed}}$$

$$-B_{\text{in}}l + \int \vec{B} \cdot d\vec{\ell} - B_{\text{out}}l + \int \vec{B} \cdot d\vec{\ell}$$

$\underbrace{\int \vec{B} \cdot d\vec{\ell}}_{=0} \quad \uparrow \quad \underbrace{\int \vec{B} \cdot d\vec{\ell}}_{=0}$
 $B_{\text{out}} = 0$

$$-NI = -nlI$$

\uparrow # of turns \uparrow turns per unit length

$$\Rightarrow -B_{\text{in}}l = -\mu_0 nlI \Rightarrow$$

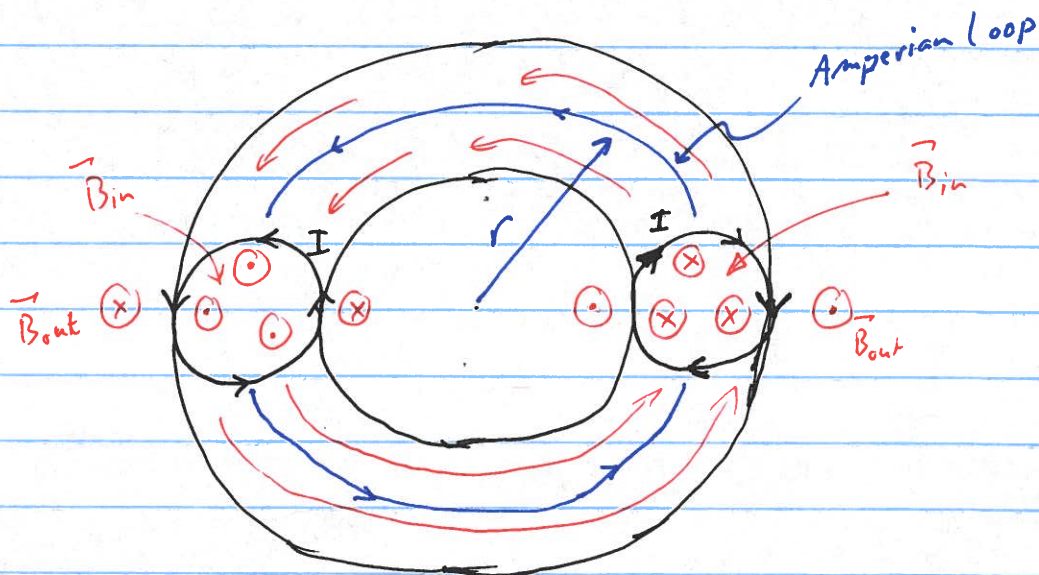
$$\vec{B}_{\text{in}} = \mu_0 n I \hat{\phi}$$

$$\vec{B}_{\text{out}} = 0$$

(no radial dependence!!)

Example: Toroidal coil [Bend solenoid into a doughnut/loop]

↳ Approximate Toroid as a series of identical current loops in the shape of a doughnut (torus).



Amperian loop outside of toroid

$$\left\{ \begin{array}{l} r < \text{toroid inner radius} \\ \text{or} \\ r > \text{toroid outer radius} \end{array} \right.$$

$$\oint_{\text{loop out}} \vec{B} \cdot d\vec{\ell} = \mu_0 \underbrace{I_{\text{enclosed}}}_{=0}$$

$$\Rightarrow -B_{\text{out}} 2\pi r = 0 \quad \Rightarrow \quad \boxed{\vec{B}_{\text{out}} = 0}$$

Amperian loop inside of toroid (toroid inner radius $< r <$ toroid outer radius)

$$\oint_{\text{loop in}} \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enclosed}}$$

$$\Rightarrow B_{\text{in}} 2\pi r = \mu_0 N I$$

total number of turns / loops

$$\Rightarrow \boxed{\vec{B}_{\text{in}} = \frac{\mu_0 N I}{2\pi r} \hat{\phi}}$$

\Rightarrow magnetic field is not constant, but it is completely enclosed inside toroid

(localized magnetic field)

Note: - Tokamak fusion "reactors" use toroidal magnets (e.g. ITER)

- Nuclei have toroidal currents due to parity violation.

↳ nuclear anapole moment.

↳ contact interaction for electrons.