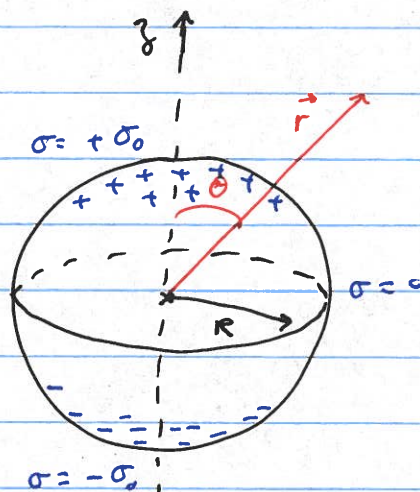


Wednesday, March 22, 2023

Separation of variables : Spherical coordinates/symmetry
(continued)

Example: Spherical shell (non-conducting) with charge density $\sigma(\theta) = \sigma_0 \cos \theta$ and radius R .

→ Calculate $V(r, \theta)$
everywhere,
(inside & outside)



direct method: $V(r, \theta) = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(\theta') r'^2 \sin \theta' dr' d\theta' d\phi'}{|\vec{r} - \vec{r}'|}$
 \downarrow somewhat messy \uparrow R^2 \uparrow no dr' integral

Separation of variables:

inside shell ($r < R$): $V_{in}(r, \theta) = \sum_{l=0}^{\infty} C_l \left(\frac{r}{R}\right)^l P_l(\cos \theta)$

outside shell ($r > R$): $V_{out}(r, \theta) = \sum_{l=0}^{\infty} C_l \left(\frac{R}{r}\right)^{l+1} P_l(\cos \theta)$

Boundary condition: At shell surface $\begin{cases} E_{\perp} \text{ is discontinuous} \\ E_{\parallel} \text{ is continuous} \end{cases}$

[see Lecture 8, Feb. 20]

$$\Delta E_{\perp} = \frac{\sigma_s}{\epsilon_0} \Rightarrow E_{\text{out},r} \Big|_{r=R} - E_{\text{in},r} \Big|_{r=R} = \frac{\sigma(\theta)}{\epsilon_0} = \frac{\sigma_0 \cos \theta}{\epsilon_0}$$

$$\left. \begin{array}{l} E_r = -\frac{\partial V}{\partial r} \\ \Leftrightarrow \end{array} \right\} \frac{\partial V_{\text{out}}}{\partial r} \Big|_{r=R} - \frac{\partial V_{\text{in}}}{\partial r} \Big|_{r=R} = -\frac{\sigma_0 \cos \theta}{\epsilon_0}$$

$$\Rightarrow \sum_{l=0}^{\infty} C_l \frac{R^{l+1}}{r^{l+2}} (-l-1) P_l(\cos \theta) \Big|_{r=R} - \sum_{l=0}^{\infty} \frac{C_l l}{R^l} r^{l-1} P_l(\cos \theta) \Big|_{r=R} = -\frac{\sigma_0 \cos \theta}{\epsilon_0}$$

$$\begin{array}{l} \downarrow r=R \\ \Rightarrow \end{array} \sum_{l=0}^{\infty} \frac{C_l (-l-1)}{R} P_l(\cos \theta) - \sum_{l=0}^{\infty} \frac{C_l l}{R} P_l(\cos \theta) = -\frac{\sigma_0 \cos \theta}{\epsilon_0}$$

$$\Leftrightarrow \sum_{l=0}^{\infty} \frac{C_l (2l+1)}{R} P_l(\cos \theta) = \frac{\sigma_0 \cos \theta}{\epsilon_0}$$

Q: How do we determine C_l 's?

A: Apply Fourier's trick! (for Legendre Polynomials)

$\hookrightarrow \begin{cases} \text{multiply both sides by } P_{l'}(\cos \theta) \\ \text{integrate both sides by } dx = d(\cos \theta) \end{cases}$

$$\Rightarrow \int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) \left[\sum_{l=0}^{\infty} \frac{C_l (2l+1)}{R} P_l(\cos \theta) \right] d(\cos \theta)$$

$$= \int_{\cos \theta = -1}^{\cos \theta = +1} \frac{\sigma_0 \cos \theta}{\epsilon_0} P_{l'}(\cos \theta) d(\cos \theta)$$

$$\Leftrightarrow \sum_{l=0}^{\infty} \frac{C_l (2l+1)}{R} \int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) P_l(\cos \theta) d(\cos \theta)$$

$$= \frac{\sigma_0}{\epsilon_0} \int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) \cos \theta d(\cos \theta)$$

$\frac{2}{2l+1} \delta_{ll'}$

Formula:

$$\int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) P_l(\cos \theta) d(\cos \theta) = \frac{2}{2l+1} \delta_{ll'}$$

or

$$\int_{x=-1}^{x=+1} P_{l'}(x) P_l(x) dx = \frac{2}{2l+1} \delta_{ll'}$$

\Rightarrow sum disappears!

$$\Rightarrow \frac{C_{l'} (2l'+1)}{R} \frac{2}{2l'+1} = \frac{\sigma_0}{\epsilon_0} \int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) \cos \theta d(\cos \theta)$$

$$\frac{2}{2l'+1} \delta_{l'1} = \frac{2}{2(1)+1} = \frac{2}{3}$$

\Rightarrow for $l' \neq 1$, $C_{l'} = 0$

$\Rightarrow l' = 1$

Note: The entire sum disappears! (this does not always happen)

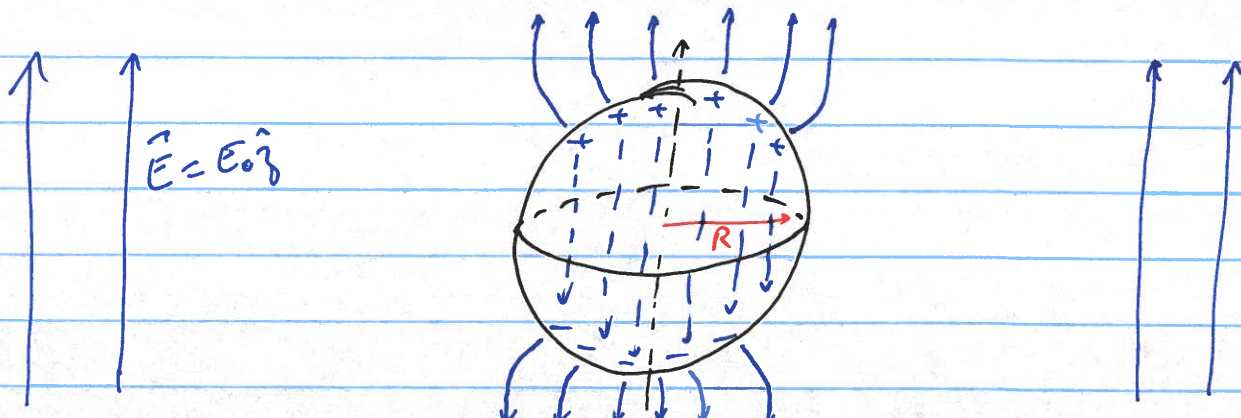
$$\text{for } l'=1, \quad \frac{C_1}{R} = \frac{\sigma_0}{\epsilon_0} \frac{2}{3}$$

$$\Rightarrow C_1 = \frac{\sigma_0 R}{\epsilon_0} \frac{2}{3}$$

$$\Rightarrow \begin{cases} V_{in} = C_1 \left(\frac{r}{R}\right)^1 P_1(\cos\theta) = \frac{\sigma_0 R}{\epsilon_0} \frac{2}{3} \frac{r}{R} \cos\theta \\ V_{out} = C_1 \left(\frac{R}{r}\right)^{1+1} P_1(\cos\theta) = \frac{\sigma_0 R}{\epsilon_0} \frac{2}{3} \frac{R^2}{r^2} \cos\theta \end{cases}$$

$$\Rightarrow \begin{cases} V_{in}(r, \theta) = \frac{\sigma_0}{3\epsilon_0} r \cos\theta = \frac{\sigma_0}{3\epsilon_0} z & (r \leq R) \\ V_{out}(r, \theta) = \frac{\sigma_0 R^3}{3\epsilon_0} \frac{\cos\theta}{r^2} & (r \geq R) \end{cases}$$

Example 2: Solid metal sphere (conducting) in an otherwise uniform electric field $\vec{E} = E_0 \hat{z}$.
 ↳ Calculate $V(r, \theta)$.



Q1: What is the potential at $|r| \rightarrow +\infty$?

A1: Not zero, Actually $V(r \rightarrow \infty) \rightarrow \pm\infty, 0!!!$
 \hookrightarrow unusual.

Q2: What is the potential inside the metal sphere?

A2: It must be an equipotential!

$$\hookrightarrow V_{\text{inside sphere}}(r, \theta) = \text{cst}$$

\uparrow
 $r < R$

For convenience, we pick $V_{\text{sphere}} = 0$ and put it at origin
 \hookrightarrow inside + surface

Q3: What is the potential far from the sphere?

A3: $\vec{E}_{\text{far}} = E_0 \hat{z} \Rightarrow V = -E_0 z + \text{cst}$
 $[\vec{E} = -\vec{\nabla} V]$ \uparrow cst = 0 by symmetry for $V_{\text{sphere}} = 0$

\Rightarrow Boundary Conditions: (i) $V = 0$ for $r \leq R$

(ii) $\vec{E}(r=R) \cdot \hat{r}$ (ii) $V \rightarrow -E_0 z$ for $r \gg R$
 $= -E_0 r \cos \theta$

General form of solution (outside, for $r \gg R$)

$$V(r, \theta) = \sum_{l=0}^{\infty} \left[a_l \left(\frac{r}{R}\right)^l + b_l \left(\frac{R}{r}\right)^{l+1} \right] P_l(\cos \theta)$$

boundary condition (i):

if $V(r=R, \theta) = 0$, then $a_l \left(\frac{R}{R}\right)^l + b_l \left(\frac{R}{R}\right)^{l+1} = 0$
 $\Rightarrow a_l = -b_l$

$$\text{Thus } V_{\text{outside}}(r, \theta) = \sum_{l=0}^{\infty} a_l \left[\left(\frac{r}{R}\right)^l - \left(\frac{R}{r}\right)^{l+1} \right] P_l(\cos \theta)$$

boundary condition (i)

$$\text{for } r \gg R, \text{ then } V(r, \theta) = \sum_{l=0}^{\infty} a_l \left[\left(\frac{r}{R}\right)^l - \left(\frac{R}{r}\right)^{l+1} \right] P_l(\cos \theta)$$

$$= -E_0 r \cos \theta$$

individual terms do not contribute at $r \gg R$

$$\text{Apply "Fourier trick": } \int_{\cos \theta = -1}^{\cos \theta = +1} P_{l'}(\cos \theta) P_l(\cos \theta) d(\cos \theta) = \frac{2}{2l+1} \delta_{l'l}$$

OR just note that $\cos \theta = P_{l=1}(\cos \theta)$

↳ only $l=1$ term is present in sum: $a_{l \neq 1} = 0$

$$\Rightarrow a_{l=1} \left(\frac{r}{R}\right) P_{l=1}(\cos \theta) = -E_0 r \cos \theta$$

$$\Rightarrow a_1 \frac{r}{R} \cos \theta = -E_0 r \cos \theta \quad \Rightarrow \quad a_1 = -E_0 R$$

Thus

$$\begin{aligned} V_{\text{outside}}(r, \theta) &= -E_0 R \left[\left(\frac{r}{R}\right) - \left(\frac{R}{r}\right)^2 \right] \cos \theta \\ &= -E_0 r \cos \theta + E_0 \frac{R^3}{r^2} \cos \theta \\ &= V_{\text{ext}} + E_0 \frac{R^3}{r^2} \cos \theta \end{aligned}$$

Q: What's the surface charge distribution on the sphere?
induced charge

$$\underline{A:} \quad -\frac{\sigma(\theta)}{\epsilon_0} = \left. \frac{\partial V_{\text{out}}}{\partial r} \right|_{r=R} - \underbrace{\left. \frac{\partial V_{\text{in}}}{\partial r} \right|_{r=R}}_{=0} = \Delta E$$

$$= -E_0 \cos \theta + E_0 R^3 \cos \theta (-2) \frac{1}{r^3} \Big|_{r=R}$$

$$= -E_0 \cos \theta (1+2)$$

$$= -3E_0 \cos \theta \Rightarrow \sigma(\theta) = 3\epsilon_0 E_0 \cos \theta$$

$$\begin{cases} \sigma(\theta < \pi/2) \text{ is positive} \\ \sigma(\theta > \pi/2) \text{ is negative} \\ \sigma(\theta = \pi/2) = 0 \end{cases}$$