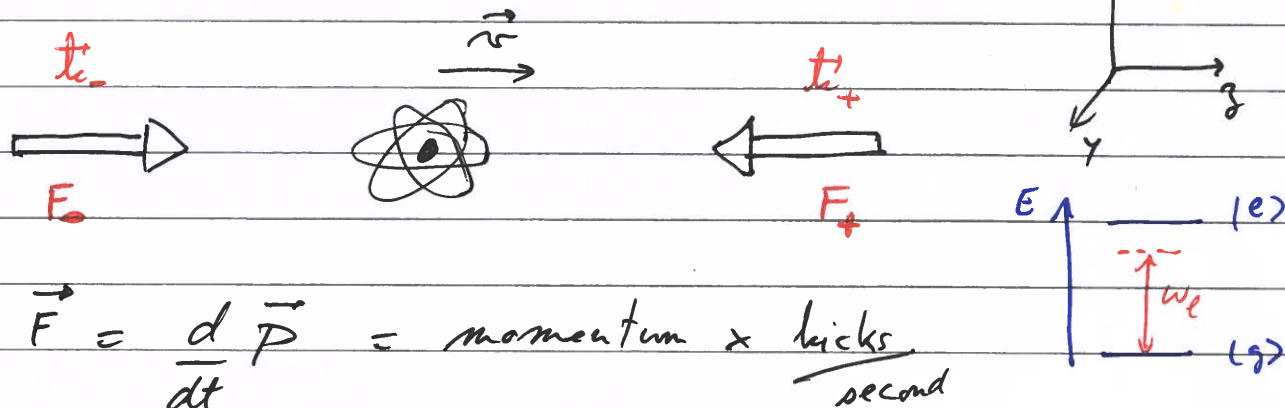


Thursday, November 13, 2025

Semi-Classical model of Doppler Cooling

$$\vec{F} = \frac{d\vec{p}}{dt} = \text{momentum} \times \frac{\text{hits}}{\text{second}}$$

$$= \hbar k \sigma_{\text{scattering}}$$

$$\rho_0 = \frac{I}{I_{\text{sat}}}$$

$$= \hbar k \frac{\rho_0}{1 + \rho_0 + \left(\frac{2\delta}{\gamma}\right)^2} \frac{\gamma}{2}$$

$$I_{\text{sat}} = 1.6 \text{ mW/cm}^2$$

for Rb

$2\pi \times 6 \text{ MHz}$
for Rb

1D model: We will do the calculation in the low intensity limit ($\rho_0 = \frac{I}{I_{\text{sat}}} \ll 1$) so that the

atoms do not have to choose between the two laser beams, i.e. saturation is not important.

detuning: $\delta = \delta_e \pm \underbrace{\Delta\omega_{\text{Doppler}}}_{\hbar \vec{k} \cdot \vec{v}} = \delta_e \pm \hbar v$

$\delta_e = \omega_e - \omega_l$

laser direction

In 1D: $F_{\text{total}} = -F_+ + F_-$

$$= \hbar k \frac{\gamma}{2} \rho_0 \left\{ \frac{-1}{1 + \rho_0 + \left[\frac{2(\delta_e + \hbar v)}{\gamma} \right]^2} + \frac{1}{1 + \rho_0 + \left[\frac{2(\delta_e - \hbar v)}{\gamma} \right]^2} \right\}$$

Low ν damping force:

Assume $|h\nu| \ll |\delta|$

$$F = \frac{t h \delta \rho_0}{2} \left\{ \frac{-1}{1 + \rho_0 + \frac{4\delta_e^2}{\gamma^2} + \frac{8h\nu\delta_e}{\gamma^2} + 4\left(\frac{h\nu}{\delta}\right)^2} + \frac{1}{1 + \rho_0 + \frac{4\delta_e^2}{\gamma^2} - \frac{8h\nu\delta_e}{\gamma^2} + 4\left(\frac{h\nu}{\delta}\right)^2} \right\}$$

neglect neglect

$$\approx \frac{t h \delta \rho_0}{2} \left\{ \frac{-1}{\left(1 + \rho_0 + \frac{4\delta_e^2}{\gamma^2}\right) \left(1 + \frac{8h\nu\delta_e/\gamma^2}{1 + \rho_0 + 4\delta_e^2/\gamma^2}\right)} + \frac{1}{\left(1 + \rho_0 + \frac{4\delta_e^2}{\gamma^2}\right) \left(1 - \frac{8h\nu\delta_e/\gamma^2}{1 + \rho_0 + 4\delta_e^2/\gamma^2}\right)} \right\}$$

$\frac{1}{1+\epsilon} \approx 1-\epsilon$ $\frac{1}{1-\epsilon} \approx 1+\epsilon$

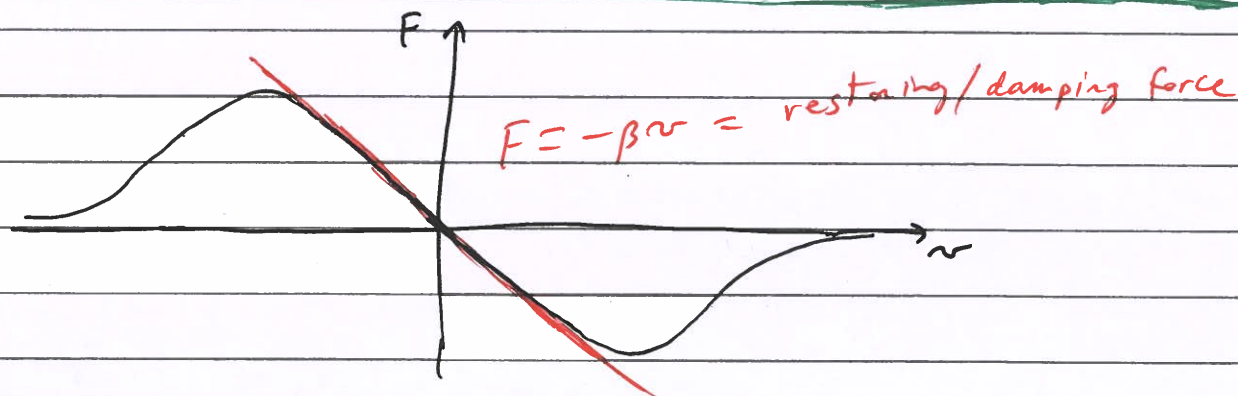
$$\approx \frac{t h \delta \rho_0}{2} \frac{1}{1 + \rho_0 + 4\delta_e^2/\gamma^2} \left\{ -1 + \frac{8h\nu\delta_e/\gamma^2}{1 + \rho_0 + 4\delta_e^2/\gamma^2} + 1 + \frac{8h\nu\delta_e/\gamma^2}{1 + \rho_0 + 4\delta_e^2/\gamma^2} \right\}$$

$$\approx \frac{t h \delta \rho_0}{(1 + \rho_0 + 4\delta_e^2/\gamma^2)^2} \frac{8h\nu\delta_e}{\gamma^2}$$

$$\Rightarrow F \approx \underbrace{\frac{8 t h^2 \delta_e \rho_0}{\gamma (1 + \rho_0 + 4\delta_e^2/\gamma^2)^2}}_{-\beta} \nu = -\beta \nu$$

damping coefficient

negative for red detuning ($\omega_L < \omega_{eg} \Rightarrow \delta_e < 0$)



Doppler Temperature

The final temperature is determined by the equilibrium condition:

$$\text{cooling rate} = \text{heating rate}$$

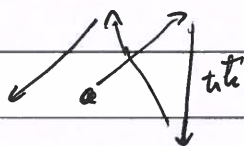
What's the cooling rate?

$$\begin{aligned} \text{Cooling rate} &= \frac{dE}{dt} = \text{power out} \\ &= \vec{F}_{\text{cooling}} \cdot \vec{v} \end{aligned}$$

$$= \frac{\rho_0}{1 + \rho_0 + \left(\frac{2\delta_L}{\gamma}\right)^2} 8\pi h^2 \delta_L v^2$$

What's the heating rate?

Random direction of absorption (3D) and random fluorescence scattering generates a random walk in momentum space with a step size of $\vec{p} = \hbar \vec{k}$
 absorption or re-emission in 3D



$$\text{heating rate} = \frac{d}{dt} \langle E_{\text{kinetic}} \rangle = \frac{d}{dt} \left[\left\langle \frac{P^2(t)}{2m} \right\rangle \right]$$

$$= \frac{1}{2m} \frac{d}{dt} \left\langle \left(\vec{p}_1 + \vec{p}_2 + \vec{p}_3 + \dots \right)^2 \right\rangle_t$$

$$= \frac{1}{2m} \frac{d}{dt} \left\langle \vec{p}_1^2 + \vec{p}_2^2 + \dots + 2\vec{p}_1 \cdot \vec{p}_2 + 2\vec{p}_2 \cdot \vec{p}_3 + \dots \right\rangle_t$$

$$= \frac{1}{2m} \frac{d}{dt} \left\{ \underbrace{\langle \vec{p}_1^2 \rangle_t}_{(\hbar k)^2} + \underbrace{\langle \vec{p}_2^2 \rangle_t}_{(\hbar k)^2} + \dots + 2\cancel{\langle \vec{p}_1 \cdot \vec{p}_2 \rangle_t} + 2\cancel{\langle \vec{p}_2 \cdot \vec{p}_3 \rangle_t} + \dots \right\}$$

$$= \frac{1}{2m} \frac{d}{dt} \left\{ \underbrace{N}_{N(t)} (\hbar k)^2 \right\}$$

no correlation between different scattering events

$$\langle \vec{p}_i \cdot \vec{p}_j \rangle_t = 0 \quad i \neq j$$

$$= \frac{1}{2m} (\hbar k)^2 \frac{d}{dt} N(t)$$

$2 \times \gamma_{\text{scattering}}$

$1+1 = \text{absorption} + \text{spontaneous emission}$

$$\Rightarrow \text{heating rate} = \frac{1}{2m} 2 (\hbar k)^2 \gamma_{\text{scattering}}$$

$$2 \times \frac{\rho_0 \gamma/2}{1 + \rho_0 + \left(\frac{2\delta_L}{\gamma}\right)^2} \quad \text{for } \langle v \rangle = 0$$

2 laser beams counter-propagating

no Doppler shift on average (neglect)

at equilibrium, heating rate = cooling rate

$$\frac{2 (\hbar k)^2}{2m} \frac{\rho_0 \gamma}{1 + \rho_0 + \left(\frac{2\delta_L}{\gamma}\right)^2} = \frac{\rho_0}{\left[1 + \rho_0 + \left(\frac{2\delta_L}{\gamma}\right)^2\right]} \frac{8\hbar^2 k^2 \delta_L v^2}{\gamma}$$

$$\Leftrightarrow \frac{2\hbar\gamma}{2m} = \frac{8\delta_L v^2}{\gamma \left[1 + \cancel{\delta_L} + \left(\frac{2\delta_L}{\gamma} \right)^2 \right]}$$

stopped here ≈ 0 for low intensity limit (neglect)

$$\Leftrightarrow \underbrace{\frac{\hbar\gamma}{2 \cdot 8} \frac{\gamma}{\delta_L} \left[1 + \left(\frac{2\delta_L}{\gamma} \right)^2 \right]} = \frac{1}{2} m v^2 = \langle E_{\text{kinetic}} \rangle$$

This function has a minimum for $\delta_L = -\gamma/2$

$$\Rightarrow \frac{\hbar\gamma}{2 \cdot 8} \frac{\gamma}{\gamma/2} \left(1 + \left(\frac{2 \cdot \cancel{\gamma/2}}{\gamma} \right)^2 \right) = \underbrace{\langle E_{\text{kinetic}} \rangle}_{\frac{1}{2} k_B T}$$

$$\Leftrightarrow \frac{\hbar\gamma}{4} = \frac{1}{2} k_B T_{\text{minimum}}$$

$$\Rightarrow T_{\text{minimum}} = \frac{\hbar\gamma}{2k_B} = \text{Doppler cooling limit}$$

or

Doppler Temperature

Ex: $T_{\text{Doppler}} \approx 180 \mu\text{K}$ for ^{87}Rb

with $\hbar = 1.054 \times 10^{-34} \text{ J}\cdot\text{s}$

$\gamma = 2\pi \times 6 \text{ MHz}$

$k_B = 1.38 \times 10^{-23} \text{ J/K}$