

Thursday, March 21, 2013

## Discrete Symmetries

### Spatial Inversion Symmetry or Parity Symmetry

Parity operation:

Classical: Parity transformation 
$$\begin{cases} x \rightarrow -x \\ y \rightarrow -y \\ z \rightarrow -z \end{cases} \quad \left( \begin{array}{l} \text{see Jackson} \\ \text{6.11 2nd ed.} \end{array} \right)$$

examples: position vector:  $\vec{r} \xrightarrow{P} -\vec{r}$  (parity odd)  
 velocity:  $\vec{v} = \frac{d\vec{r}}{dt} \xrightarrow{P} -\frac{d\vec{r}}{dt} = -\vec{v}$  (odd)

momentum:  $\vec{p} = m\vec{v} \xrightarrow{P} -m\vec{v} = -\vec{p}$  (odd)

angular momentum:  $\vec{L} = \vec{r} \times \vec{p} \xrightarrow{P} (-\vec{r}) \times (-\vec{p}) = \vec{L}$  (even)

Force:  $\vec{F} = m\vec{a} = m \frac{d^2\vec{r}}{dt^2} \xrightarrow{P} m \left( -\frac{d^2\vec{r}}{dt^2} \right) = -\vec{F}$  (odd)

Energy:  $E = \frac{1}{2}m\vec{v}^2 \xrightarrow{P} \frac{1}{2}m(-\vec{v})^2 = E$  (even)

Electric & magnetic fields: consider Lorentz force law

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$$

$$\Rightarrow \begin{cases} \vec{E} \xrightarrow{P} -\vec{E} & \text{odd} \\ \vec{B} \xrightarrow{P} \vec{B} & \text{even} \end{cases}$$

Terminology: vector quantities that are ~~odd~~ under parity are called "vectors" (or Polar vectors)  
 even under parity are called "pseudovectors"

note: you cannot accomplish a parity transformation through a series of rotations.

Quantum Mechanics: Parity (transformation) operator  $\Pi$

$$\Pi |\psi\rangle = ?$$

$$\text{we require } \langle \psi | \Pi^\dagger X \Pi | \psi \rangle = - \langle \psi | X | \psi \rangle$$

$$\Rightarrow \Pi^\dagger X \Pi = -X$$

$\Rightarrow \Pi$  should be unitary since ~~norm~~ norm should not change (or ~~probabilities~~ probabilities)

$$\langle \psi | \Pi^\dagger (\Pi |\psi\rangle) = \langle \psi | \psi \rangle$$

$$\Rightarrow \Pi^\dagger \Pi = \mathbb{1} \quad \Leftrightarrow \Pi^\dagger = \Pi^{-1}$$

$$\text{thus } \Pi \Pi^\dagger X \Pi = -\Pi X$$

$$\Rightarrow X \Pi = -\Pi X \quad \Rightarrow \Pi X + X \Pi = 0$$

Consider  $|\psi\rangle = |x\rangle$  (position state)

$$X \Pi |x\rangle = -\Pi X |x\rangle = -\Pi x |x\rangle = (-x) \Pi |x\rangle$$

$$\text{thus } X(\Pi|x\rangle) = (-x)(\Pi|x\rangle)$$

$$\Rightarrow \Pi |x\rangle = e^{i\delta} |-x\rangle \equiv |-x\rangle$$

by convention ( $\delta=0$ )

$$\boxed{\Pi |x\rangle = |-x\rangle}$$

$$\text{note } \Pi(\Pi|x\rangle) = \Pi|-x\rangle = |x\rangle \Rightarrow \Pi^2 = \mathbb{1}$$

$$\Rightarrow \Pi = \Pi^{-1} = \Pi^\dagger$$

$\Rightarrow \Pi$  is Hermitian

the eigenvalues of  $\Pi$  are  $\pm 1$ .

~~the eigenvalues of  $\Pi$  are  $\pm 1$~~

consider  $|\psi\rangle = |p\rangle$

$$\text{we expect } \Pi^\dagger P \Pi = -P \Rightarrow P \Pi = -\Pi P$$

by same argument  $\Pi |p\rangle = \pm | -p \rangle$   
or for  $|x\rangle$

How do wavefunctions transform under  $\Pi$ ?

$$\psi(\vec{r}) = \langle \vec{r} | \psi \rangle$$

$$\Pi |\psi\rangle \rightarrow \langle \vec{r} | \Pi |\psi\rangle = \langle -\vec{r} | \psi \rangle = \psi(-\vec{r})$$

If  $|\psi\rangle$  is an eigenstate of  $\Pi$ , then

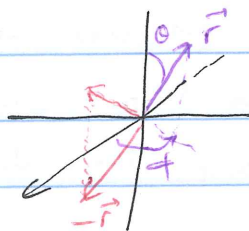
$$\Pi |\psi\rangle = \pm |\psi\rangle \text{ then } \psi(-\vec{r}) = \begin{cases} \psi(\vec{r}) & \text{even parity} \\ -\psi(\vec{r}) & \text{odd parity} \end{cases}$$

Spherical Harmonics under  $\Pi$

$$\begin{cases} x \rightarrow -x \\ y \rightarrow -y \\ z \rightarrow -z \end{cases}$$

$(\Rightarrow)$

$$\begin{cases} r = \sqrt{x^2 + y^2 + z^2} \rightarrow r \\ \theta \rightarrow \pi - \theta \\ \phi \rightarrow \phi + \pi \end{cases}$$



Since

$$Y_l^m = (-1)^m \frac{(2l+1)(l-m)!}{4\pi(l+m)!} P_l^m(\cos\theta) e^{im\phi} \quad (\text{for } m \geq 0)$$

note:  $\cos\theta \xrightarrow{P} \cos(\pi-\theta) = -\cos\theta$

$$e^{im\phi} \xrightarrow{P} e^{im(\phi+\pi)} = e^{im\pi} e^{im\phi} = (-1)^m e^{im\phi}$$

one can show that

$$Y_l^m(\theta, \phi) \xrightarrow{P} Y_l^m(\pi-\theta, \phi+\pi) = (-1)^l Y_l^m(\theta, \phi)$$

thus  $|Y_l^m\rangle$  are eigenstates of  $\Pi$ .

Hamiltonian under  $\Pi$

Most Hamiltonians are even under a parity transformation

for example  $H = \frac{P^2}{2m} + \frac{e^2}{R}$  is even

In fact  $[H, \Pi] = 0$

Thus we can construct an eigenbasis of  $H$  in which the eigenstates are odd and even functions.

$(n, l, m)$  have a definite parity given by  $l$ .  $(\alpha|nS\rangle + \beta|nP\rangle$  is an eigenstate of  $H$  but does not have a definite parity)

Parity violation:

All known ~~in~~ fundamental interactions conserve parity, except the weak interaction.

gravity, E & M, strong force

All the Hamiltonians that we have studied so far are even under parity (i.e. Hydrogen).

$$H_p = \frac{\overset{\text{even}}{P^2}}{2m} + \frac{e^2}{\underset{\text{even}}{R}} + \frac{\overset{\text{even}}{P^4}}{8m_e^3 c^2} + \frac{1}{2} \frac{e^2}{m_e^2 c^2} \frac{1}{\underset{\text{even}}{R^3}} \vec{L} \cdot \vec{S}$$

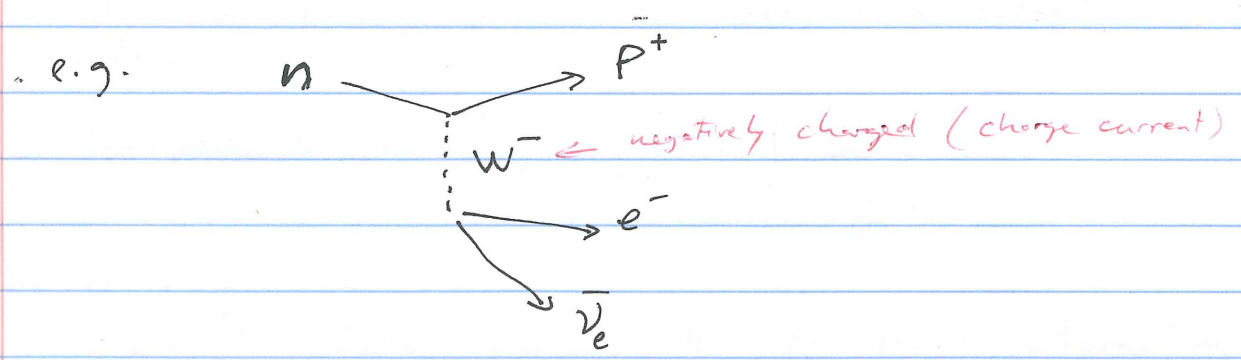
$$+ \frac{A}{\hbar^2} \vec{I} \cdot \vec{J} - \frac{q_e}{2m_e} (\vec{L} + 2\vec{S}) \cdot \vec{B} - q_e \vec{E} \cdot \vec{R} + \dots$$

↑ ↑
↑ ↑ ↑
↑ ↑

even even
even even even
odd odd

Stark shift appears to violate parity, but only if you don't transform E

The Weak interaction is responsible for radioactive decay

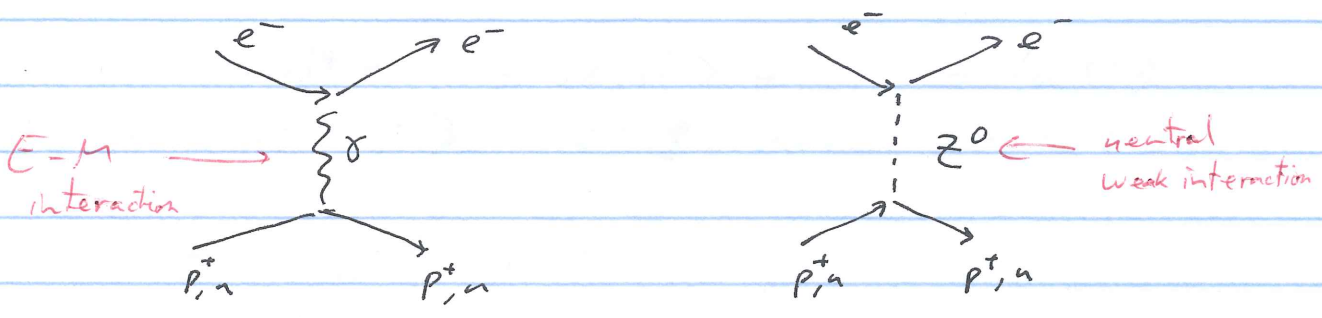


$\beta^-$  decay in  $^{60}\text{Co}$  shows an ~~asymmetry~~ asymmetry in the emission of  $e^-$

↑ ↓ ↓ ↓

I of  $^{60}\text{Co}$  (Yang, Lee, Wu 1956-57)

There is also a <sup>weak</sup> neutral current interaction mediated by  $Z^0$



In the non-relativistic limit, this interaction can be written as

$$H_{\text{int}} = \frac{G_F}{\sqrt{2}} \frac{1}{2m_e c^2} Q_w \vec{S} \cdot \left[ \vec{P} \delta^3(\vec{R}) + \delta^3(\vec{R}) \vec{P} \right]$$

$\vec{e} \text{ spin}$      $\vec{e} \text{ momentum}$      $\vec{e} \text{ position}$   
 ↑    ↑    ↑    ↑    ↑  
 even    odd    even    even    odd  
 —————  
 odd

$$G_F \approx 3 \times 10^{-12} m_e c^2 \left( \frac{t}{m_e c} \right)^3$$

$$\text{and } Q_w = -N + (1 - 4 \sin^2 \theta_w) Z$$

$$\approx -N$$

= Weak charge of nucleus

~~$\vec{S} \cdot \vec{P}$~~  is odd under parity  $\rightarrow$  pseudoscalar

In the standard model  $\sin^2 \theta_w \approx 0.23$

( $\theta_w$  is called the Weinberg mixing angle)

$H_{Z_0}$  is parity odd but very very small.

~~In principle~~ Thus  $[H_0 + H_{Z_0}, \pi] \neq 0$   
 $\approx 0$

$\delta^3(\vec{R})$  term indicates that the weak interaction is a contact interaction and only affects the  $e^-$  inside the nucleus.